# **Nano-Micro Letters**

# ARTICLE

https://doi.org/10.1007/s40820-025-01776-3



Cite as Nano-Micro Lett. (2025) 17:270

Received: 31 January 2025 Accepted: 17 April 2025 © The Author(s) 2025

# Efficient Thermally Evaporated Near-Infrared Perovskite Light-Emitting Diodes via Phase Regulation

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# HIGHLIGHTS

- α-phase formamidinium lead triiodide (FAPbI<sub>3</sub>) was prepared based on triple-source co-evaporation.
- A partial Cs-doping in the FAPbI<sub>3</sub> can help with the suppression of the non-radiative recombination, elimination of the metallic Pb, improvement of spatial confinement.
- Near-infrared perovskite light-emitting diodes (NIR-PeLEDs) based on triple-source co-evaporated FACsPbI<sub>3</sub> thin films achieved a maximum external quantum efficiency of 10.25%, which was around 6 times higher than that of FAPbI<sub>3</sub>-based NIR-PeLEDs.

**ABSTRACT**  $\alpha$ -phase formamidinium lead triiodide (FAPbI<sub>3</sub>) has demonstrated extraordinary properties for near-infrared perovskite lightemitting diodes (NIR-PeLEDs). The vacuum processing technique has recently received increasing attention from industry and academia due to its solvent-free feature and compatibility with large-scale production. Nevertheless, vacuum-deposited NIR-PeLEDs have been less studied, and their efficiencies lag far behind those of solution-based PeLEDs as it is still challenging to prepare pure  $\alpha$ -FAPbI<sub>3</sub> by the thermal evaporation. Herein, we report a Cs-containing triple-source co-evaporation approach



to develop the perovskite films. The addition of thermally stable Cs cation fills in the perovskite crystal lattice and eliminates the formation of metallic Pb caused by the degradation of FA cation during the evaporation process. The tri-source co-evaporation strategy significantly promotes the phase transition from yellow  $\delta$ -phase FAPbI<sub>3</sub> to black  $\alpha$ -phase FACsPbI<sub>3</sub>, fostering smooth, uniform, and pinhole-free perovskite films with higher crystallinity and fewer defects. On this basis, the resulting NIR-PeLED based on FACsPbI<sub>3</sub> yields a maximum EQE of 10.25%, which is around sixfold higher than that of FAPbI<sub>3</sub>-based PeLEDs. Our work demonstrates a reliable and effective strategy to achieve  $\alpha$ -FAPbI<sub>3</sub> via thermal evaporation and paves the pathway toward highly efficient perovskite optoelectronic devices for future commercialization.

**KEYWORDS** Vacuum-deposited perovskite; Near-infrared emission; Phase regulation; Co-evaporation; Electroluminescence

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## **1** Introduction

Metal halide perovskites have rapidly emerged as a revolutionary frontier in display and lighting owing to their fascinating properties including high color purity, excellent charge transport properties, tunable chromaticity, and cost-effectiveness [1–5]. Specially, near-infrared perovskite light-emitting diodes (NIR-PeLEDs) hold great potential for a variety of applications such as medical treatment, data storage, optical communications, hyperspectral imaging, night vision for surveillance, and automotive safety [6, 7].

With tremendous efforts of researchers, the external quantum efficiencies (EOEs) of NIR-PeLEDs have exceeded 23% within the last few years [2, 4], enabling them to be a strong competitor of the III-V NIR-LEDs. Up to date, the fabrication of formamidinium-based emitters such as FAPbI<sub>3</sub> was regarded as the most successful approach to develop NIR-PeLEDs [2, 4, 7–9]. Nevertheless, most of the reported state-of-the-art NIR-PeLEDs were prepared via solutionbased spin-coating, limiting their application in large-scale production. Inspired by the deposition of commercialized organic light-emitting diodes (OLEDs), thermal evaporation has been proposed as an alternative and effective technique to develop scalable PeLEDs [10]. Unlike solution-processing techniques, thermal evaporation can effectively avoid the problem of low solubility of some perovskite precursors. Besides, it can prevent the utilization of environmentally unfriendly organic solvents such as dimethylformamide (DMF) and dimethyl sulfoxide (DMSO). Furthermore, the process of thermal evaporation can be uniformly fabricated onto various substrates, paving a novel pathway for the heterogeneous integration of NIR-PeLEDs. Most importantly, thermal evaporation is compatible with the currently available OLED industry, thus lowering the initial investment and speeding up the commercialization of PeLEDs [3, 10–13].

The present researches on thermally evaporated PeLEDs are mainly focused on the green emission field [3, 11, 14–18]. The studies on thermally evaporated NIR-PeLEDs have not yet been investigated in-depth. Meanwhile, the efficiencies of thermally evaporated NIR-PeLEDs are still lagging far behind those of solution-processed NIR-PeLEDs [19]. Even though FAPbI<sub>3</sub> has been widely reported in the solution-processed NIR-PeLEDs, it is believed that the direct reaction between formamidinium iodide (FAI) and lead iodide (PbI<sub>2</sub>) during the vapor deposition process usually causes

the  $\delta$ -phase and  $\alpha$ -phase mixed FAPbI<sub>3</sub> [20, 21]. This issue was demonstrated to be associated with the degradation of FA cation with the elevated evaporation temperature [22]. It is well known that the presence of hexagonal  $\delta$ -phase limits the use of pure FAPbI<sub>3</sub> as the active layer in both perovskite solar cells and PeLED [23-25]. The incorporation of mixedcation (FA/MA/Cs) and/or mixed-halide (I/Br/Cl) has been frequently implemented into converting vellow  $\delta$ -phase nonperovskite to a black  $\alpha$ -phase in the solution-prepared FAPbI<sub>3</sub> [26–29]. Compared with single cation FA-based perovskite, binary A-site cation FA/Cs-based perovskite has led to more efficient and stable perovskite due to the higher thermal stability of Cs cation [30]. Nevertheless, the employment of mixed-cation and mixed-halide stoichiometries in the thermal evaporation is technically difficult because of the complication of fabricating multiple perovskite precursors and additives simultaneously [21].

Herein, we proposed a partial replacement of FA cation by Cs cation to stabilize the perovskite crystal structure and prepare  $\alpha$ -phase FACsPbI<sub>3</sub> thin film by simultaneously thermal co-evaporation of three source, *i.e.*, FAI, PbI<sub>2</sub>, and CsI, for the application of NIR-PeLEDs. Subsequently, we systematically studied how the triple-source co-evaporation forms pure black  $\alpha$ -phase FACsPbI<sub>3</sub> in the vacuum deposition process. We demonstrated that alloying partial Cs into FAPbI<sub>3</sub> can lead to not only the enhancement of the  $\alpha$ -phase crystallinity but also the improvement of surface morphologies. Meanwhile, the FACsPbI<sub>3</sub> thin films exhibited reduced nanocrystal size and suppressed trap densities, resulting in a significant enhancement of spatial confinement and radiative recombination. Leveraging these benefits, we fabricated the NIR-PeLEDs based on thermally evaporated FACsPbI<sub>3</sub> thin films and achieved a maximum EQE value of 10.25%, which was around six times higher than that of FAPbI<sub>3</sub>-based NIR-PeLEDs.

#### **2** Experimental Section

#### 2.1 Materials

Formamidinium iodide (FAI, > 99.99%) was purchased from GreatCell Solar. Lead iodide (PbI<sub>2</sub>, 99.999%), cesium iodide (CsI, 99.999%), lithium fluoride (LiF, > 99.99%), chlorobenzene (99.8%, anhydrous) were purchased from Sigma-Aldrich. Poly(ethylene dioxythiophene)/polystyrene sulfonate (PEDOT:PSS, AI 4083) was purchased from Xi'an Yuri Solar Co. Ltd. Poly(9,9-dioctylfluoreneyl-2,7-diyl)alt-(4,4'-(N-(4-butylphenyl)-diphenylamine) (TFB) and 1,3,5-tris(1-phenyl-1H-benzimidazol-2-yl)benzene (TPBi) were purchased from Xi'an Polymer Light Technology Corp.

#### 2.2 Preparation of Perovskite Light-Emitting Diodes

#### 2.2.1 Preparation of Perovskite Films

For the perovskite deposition, a thermal evaporation integrated into an N<sub>2</sub> glovebox was used. The system was firstly vacuumed to a base pressure of  $1.0 \times 10^{-5}$  Torr before the source was heated. Three quartz crystal microbalances (QCMs) were employed to measure the rate of each source independently. Prior to deposition, the deposited rate and thickness of PbI<sub>2</sub> and FAI were calibrated using QCM and Alpha-step profiler. For the evaporation of FAPbI<sub>3</sub>, FAI and PbI<sub>2</sub> were kept in the independent crucibles and slowly heated to a target rate of 0.1 and 0.05 Å s<sup>-1</sup>, separately. For the evaporation of FACsPbI<sub>3</sub>, the evaporated rates of FAI, PbI<sub>2</sub>, and CsI were maintained at 0.1, 0.05, and 0.02 Å s<sup>-1</sup>, respectively. Due to the complicated evaporated behavior of FAI, we kept the FAI source at 55 °C for 30 min before gradually heating to the sublimation temperature (~100-120 °C for FAI) to avoid abrupt evaporation and decomposition. At the same time, the vacuum pressure of the system was increased to ~ $5.0 \times 10^{-4}$  Torr. It should be noted that the substrates holder did not rotate during the evaporation.

#### 2.2.2 Preparation of Perovskite Light-Emitting Diodes

The patterned indium tin oxide (ITO) glass substrates were sequentially sonicated in ethanol and water for 20 min. Afterward, the substrates were treated with UV-ozone for 15 min. The PEDOT:PSS was spin-coated onto the precleaned ITO-patterned glass substrate at 2000 r min<sup>-1</sup> for 30 s followed by annealing at 150 °C for 30 min. The PEDOT:PSS coated substrates were then transferred to the N<sub>2</sub> glovebox. In the glovebox, a TFB solution dissolved in chlorobenzene at a concentration of 5 mg mL<sup>-1</sup> was spin-coated onto the PEDOT:PSS film at 3000 r min<sup>-1</sup> for 30 s. The TFB layer was baked at 150 °C for 30 min. After that, these samples were transferred to a thermal evaporation

chamber integrated with a glovebox for the deposition of perovskite emitting layer. The emitting layer of FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> was deposited by co-evaporation of FAI, PbI<sub>2</sub>, and CsI in individual crucibles. The thickness of emission layer was around 20 nm. After that, TPBi (40 nm), LiF (1 nm), and aluminum (90 nm) were thermally evaporated onto the perovskite films.

#### 2.3 Characterizations

#### 2.3.1 Characterization of Perovskite Films

X-ray diffraction (XRD) spectra of perovskite films were obtained by Smart Lab 9KW multifunctional rotating-anode X-ray diffractometer using Cu Kα radiation. The absorption spectra of the perovskite films were analyzed using an ultraviolet-visible (UV-Vis-near-IR) spectrophotometer (Shimadzu Instruments, SolidSpec-3700). The steadystate photoluminescence (PL) spectra were tested using a fluorescence spectrometer (Horiba, QuantaMaster 8000). Time-resolved photoluminescence (TRPL) decay was measured using a Horiba Fluorolog-QM series. A single-mode 440 nm pulsed diode laser (pulse width of 75 ps; average power: 2.4 mW) was used as an excitation source. X-ray photoelectron spectroscopy (XPS) and ultraviolet photoelectron spectroscopy (UPS) of perovskite films were analyzed using X-ray photoelectron spectrometer (AXIS-ULTRA DLD-600W, Shimadzu) system. The transient absorption (TA) measurement with similar pump fluences was conducted by a commercial equipment (Ultrafast system, Helios fire) with a 400 nm wavelength pulse laser. The topography and surface morphology were characterized with an atomic force microscope (AFM, Jupiter XR, Oxford Instruments Asylum Research, Inc.) and field-emission scanning electron microscope (FE-SEM, FEI Nova NanoSEM 450).

# 2.3.2 Characterization of Perovskite Light-Emitting Diodes

The active area of PeLED devices was 4 mm<sup>2</sup> which was defined by the overlapping area of ITO and Al electrodes. The electroluminescence (EL) performance including current density–voltage–radiance characteristics, external quantum efficiency, and EL spectra of PeLEDs were collected

by a computer-controlled integrated system, consisting of a programmable Keithley 2400 source meter and XPQY-EQE (Guangzhou XiPu Optoelectronics Technology) equipped with an integrating sphere and a photodetector. The measurement system was calibrated by halogen lamps metered by the National Institute of Standards and Technology (NIST) and installed inside an  $N_2$ -filled glovebox.

# **3** Results and Discussion

#### 3.1 Preparation of Thermally Evaporated Perovskites

In this work, we deposited the dual-source evaporated FAPbI<sub>3</sub> and triple-source evaporated FACsPbI<sub>3</sub> perovskite thin films on a steady substrate to develop spatially gradient components according to previously reported publications [11, 18]. Utilizing a fixed substrate simplifies maintaining stoichiometric consistency, which is critical for high-throughput optimization process. However, substrate rotation introduces challenges in maintaining precious compositional control. In Fig. S1, the as-deposited spatially gradient FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> thin films obtained via thermal evaporation were illustrated. There was a gradient color distribution in the FAPbI<sub>3</sub> thin films, and most FAPbI<sub>3</sub> thin films exhibited a yellowish color, while all the FACsPbI<sub>3</sub> thin films showed a brown color. Unlike the easy deposition of organic materials in OLEDs, the evaporated process of perovskite is complicated, involving precursor evaporation accompany with decomposition, chemical reaction, nucleus formation, and crystal growth. Figure 1a illustrates the preparation process of FAPbI<sub>3</sub> thin films based on the direct reaction between FAI and PbI<sub>2</sub> using the thermal evaporation process. During the evaporation of FAI, it can be easily degraded to hydrogen cyanide and hydrogen iodide, which further react to form 1.3.5-sym-triazine and ammonia. With increasing temperature, FA may degrade and leave the crystal lattice, remaining the unreacted PbI2 and metallic Pb. This leads to a yellowish color in the as-deposited FAPbI<sub>3</sub> thin film, indicating the presence of  $\delta$ -phase FAPbI<sub>3</sub> (Fig. 1a). Typically, the  $\delta$ -phase FAPbI<sub>3</sub> has a hexagonal structure that leads to non-photoactivity [20, 31]. To address this problem, co-evaporation of FAI, PbI<sub>2</sub>, and CsI was proposed to convert the yellow  $\delta$ -phase to the black  $\alpha$ -phase and stabilize the perovskite phase. During the evaporation, the thermally stable Cs cation can fill the empty A-cation sites

and stabilize the crystal structure of FAPbI<sub>3</sub> (Fig. 1b). The photograph of the as-deposited FACsPbI<sub>3</sub> thin film exhibited a different brown color from that of FAPbI<sub>3</sub> thin film. This can be ascribed to the dominant  $\alpha$ -FAPbI<sub>3</sub> phase in the FACsPbI<sub>3</sub> thin film, which has an ordered cubic structure.

## 3.2 Fundamental Photoelectronic Properties of Perovskites

To understand the crystalline phases formed in the FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> thin films, we conducted X-ray diffraction (XRD) measurements. As shown in Fig. 2a, the diffraction peaks of the as-evaporated FAPbI<sub>3</sub> at 11.7° and 14.0° evidenced the coexistence of  $\delta$ -phase and  $\alpha$ -phase FAPbI<sub>2</sub> [23, 26, 32]. Notably, the  $\delta$ -FAPbI<sub>3</sub> phase at 11.7° was almost absent in the FACsPbI<sub>3</sub> thin film, confirming the conversion to the  $\alpha$ -FAPbI<sub>3</sub> phase after co-evaporation with CsI. Furthermore, Fig. 2b compares the UV-Vis absorption spectra of FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> thin films. Both FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> thin films had a characteristic absorption edge at around 800 nm, ascribed to the absorption edge of the  $\alpha$ -phase FAPbI<sub>3</sub> [23]. It was noted that there was a sharp absorption peak at 380 nm which as attributed to the  $\delta$ -FAPbI<sub>3</sub> phase [21]. However, no obvious  $\delta$ -phase signal was observed in the FACsPbI<sub>3</sub> thin films. The photoluminescence (PL) spectra of the co-evaporated FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> are also shown in Fig. 2b. The PL peak of the FAPbI<sub>3</sub> thin film was located at a wavelength of 784 nm with a full-width at half-maximum (FWHM) of 55 nm. In contrast, the PL peak of the FACsPbI<sub>3</sub> thin film was at 782 nm with a narrower FWHM of 45 nm. The narrower PL emission observed in the triple-source evaporated FACsPbI<sub>3</sub> thin films may be an indication of improved monodispersity and reduced defects [33]. In addition, the PL quantum yield (PLQY) of the FACsPbI<sub>3</sub> thin film reached ~ 26%, which is more than eight times as high as that of the FAPbI<sub>3</sub> thin film (PLQY  $\sim 3\%$ ). The PL intensity of the FACsPbI<sub>3</sub> thin film was much higher than that of the FAPbI<sub>3</sub> thin film, suggesting an enhanced PL emission and reduced non-radiative recombination in the FACsPbI<sub>3</sub> thin film. To confirm this, time-resolved PL (TRPL) measurement was performed to assess the carrier recombination dynamics of FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> thin films. The TRPL spectra were fitted by bi-exponential decay (Eq. S1), and the average PL lifetime was determined using Eq. S2. The



Fig. 1 Preparation of thermally evaporated perovskite thin films. a Dual-source co-evaporation of  $FAPbI_3$  formed via FAI and  $PbI_2$ , and b triple-source co-evaporation of  $FACsPbI_3$  formed via FAI, CsI, and  $PbI_2$ 

fitting details are summarized in Table S1. As illustrated in Fig. 2c, FACsPbI<sub>3</sub> thin film showed a longer average carrier lifetime of 51.02 ns than that of FAPbI<sub>3</sub> thin film with an average lifetime of 48.55 ns, revealing the reduced defects and non-radiative recombination. XPS measurements were further conducted to investigate the detailed chemical states of the FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> thin films (Fig. S2). In Fig. 2d, there were two main peaks Pb  $4f_{5/2}$ and Pb  $4f_{7/2}$  at 143.1 and 138.3 eV, respectively. Notably, the doublet peaks at around 141.3 and 136.4 eV corresponding to the metallic Pb states were obvious in the FAPbI<sub>3</sub> thin film. The presence of a large substantial of atomic Pb was likely to be Pb-interstitial defects in the perovskite lattice of FAPbI<sub>3</sub> [34, 35]. The atomic Pb species can serve as trap states for the non-radiative recombination and would greatly undermine the optoelectronic properties of PeLEDs. In contrast, the metallic Pb state peaks were not observed in the triple-source evaporated FACsPbI<sub>3</sub>. As shown in Fig. 2e, no Cs 3d peaks were detected from FAPbI<sub>3</sub> thin film, while they were observed from FACsPbI<sub>3</sub> thin film, confirming the successful addition of Cs via triple-source co-evaporation. Furthermore,

we calculated the ratio of Cs/Pb via the integrated areas of XPS spectra for Pb 4*f* and Cs 3*d* peaks. The ratio of Cs/ Pb was determined to be 12% for the optimized FACsPbI<sub>3</sub> thin film. In the I 3*d* core-level spectra, both FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> exhibited two dominant I  $3d_{3/2}$  and I  $3d_{5/2}$  peaks (Fig. 2f). Compared with FAPbI<sub>3</sub>, the I  $3d_{3/2}$  and I  $3d_{5/2}$ peaks of FACsPbI<sub>3</sub> slightly shift toward higher binding energy, which may be ascribed to the stronger bond energies between Pb and I [36].

The energy levels of FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> were estimated by ultraviolet photoelectron spectroscopy (UPS), as illustrated in Fig. S3. The cutoff energy ( $E_{cutoff}$ ) and onset energy ( $E_{onset}$ ) of FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> are shown in Fig. 2g, h. The optical bandgap ( $E_g$ ) values of FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> thin films were fitted from the Tauc plots in Fig. 2i. The  $E_g$  values of the FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> were 1.60 and 1.54 eV, respectively. The valance band maximum ( $E_{VBM}$ ), Fermi level ( $E_{Fermi}$ ), and conduction band minimum ( $E_{CBM}$ ) values of FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> were calculated using Eqs. S3-S5 and are summarized in Table S2. The  $E_{CBM}$  values of FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> were determined to be -5.69 and -5.66 eV, respectively. Compared with FAPbI<sub>3</sub>



**Fig. 2 a** XRD patterns, **b** UV–Vis absorption and PL spectra, **c** time-resolved PL spectra and corresponding fitting spectra of FAPbI<sub>3</sub> and FACsPbI<sub>3</sub>. XPS spectra of the **d** Pb 4*f*, **e** Cs 3*d*, and **f** I 3*d*, UPS spectra at the **g** cutoff region, and **h** onset region, **i** Tauc plot of FAPbI<sub>3</sub> and FACsPbI<sub>3</sub>

thin films, there was a small upshift of  $E_{\text{CBM}}$  in the FACsPbI<sub>3</sub> thin films, leading to better hole transportation.

Transient absorption (TA) spectroscopy was applied to study the carrier dynamics in the FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> thin films. The two-dimensional TA spectra as a function of wavelength and pump–probe delay time for the FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> thin films are illustrated in Fig. 3a, b. The TA contour maps of both the FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> thin films exhibited a distinct ground-state bleaching (GSB) peak at around 770 nm and the photoabsorption (PA) positive band at 520–700 nm, which are in agreement with earlier publications [2, 37]. By comparing the GSB peak at around 770 nm between FAPbI<sub>3</sub> and FACsPbI<sub>3</sub>, it was noted that the negative peak intensity of the FACsPbI<sub>3</sub> was greatly enhanced. For better comparison, TA spectra at selected timescales of the FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> thin films are extracted in Fig. 3c, d. Once the perovskite thin films absorb photons, the photoexcited carriers owing to the bandgap renormalization effect would relax to the lowest energy sites, resulting in a redshift of the GSB [38–40]. After a photoexcitation time of 100 ps, it was obvious that the FAPbI<sub>3</sub> showed a greater redshift of ~5 nm than that of FACsPbI<sub>3</sub> (redshift: ~2 nm). The reduced redshift of FACsPbI<sub>3</sub> thin film suggested a flatter energy landscape and reduced tail states below the bandgap. Furthermore, as shown in Fig. 3e, f, TA traces as a function of delay time

at a probing wavelength of ~770 nm were extracted and fitted with a tri-exponential function. The fast decay  $(\tau_1)$  is attributed to the relaxation of photogenerated hot carriers. A higher carrier concentration can slow hot carrier cooling due to screened electron-phonon interactions [41]. The intermediate decay  $(\tau_2)$  corresponds to the carrier trapping of defects. Compared with FAPbI<sub>3</sub>, FACsPbI<sub>3</sub> showed a prolonged  $\tau_2$ , suggesting reduced carrier quenching sites. Consequently, the bleach signal for the FACsPbI<sub>3</sub> thin film exhibited a higher fraction of long-lived component (longest decay  $\tau_3 = 4345$  ps) than that of the FAPbI<sub>3</sub> thin film (longest decay  $\tau_3 = 2668$  ps), which was responsible for the recombination of photoexcited species [40]. The higher fraction of the long-lived component  $(\tau_3)$  in FACsPbI<sub>3</sub> suggests a greater contribution from radiative recombination, which is also consistent with the enhanced photoluminescence intensity, leading to higher efficiency and better luminescence properties [42].

#### 3.3 Morphology Studies of Perovskites

To investigate the effects of Cs addition to the  $FAPbI_3$  on the topographies and surface morphologies, we conducted atomic force microscope (AFM) and field-emission scanning electron microscope (FE-SEM) measurements of evaporated FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> on indium tin oxide (ITO) substrate. As shown in Fig. 4a, b, the FACsPbI<sub>3</sub> exhibited a lower surface roughness with a root-mean-square (RMS) value of 5.35 nm compared with that of FAPbI<sub>2</sub> at 7.61 nm. The lower roughness of FACsPbI<sub>3</sub> could form better contact with the hole transport layer, enabling better hole injection and transportation. Furthermore, the FACsPbI<sub>3</sub> perovskite film had a smaller nanograin compared with that of FAPbI<sub>3</sub> (Fig. 4c, d). As illustrated in Fig. 4e, f, the resulting evaporated FACsPbI<sub>3</sub> film had a smaller nanocrystal size of ~ 61.6 nm than that of FAPbI<sub>3</sub> with an average nanocrystal size of ~102.9 nm. For better comparison, we also deposited the perovskite thin films on the PEDOT:PSS/ TFB hole transport layers (HTLs). As illustrated in Fig. S4, the resulting evaporated FACsPbI3 exhibited much more compact and dense morphology compared with FACsPbI<sub>3</sub>. Besides, FACsPbI<sub>3</sub> film still had a smaller nanocrystal size of ~77.7 nm than that of FAPbI<sub>3</sub> with an average nanocrystal size of ~135.6 nm. The partial substitution of smaller Cs may eliminate the large crystallite of  $\delta$ -phase FAPbI<sub>3</sub> and suppress the smaller remnant inclusions that acted as the potential defects, leading to intrinsic spatial confinement in FACsPbI<sub>3</sub> [29]. The enhanced spatial confinement effect of



Fig. 3 Time-wavelength-dependent TA contour maps of a FAPbI<sub>3</sub> and b FACsPbI<sub>3</sub>. TA spectra at selected timescales of c FAPbI<sub>3</sub> and d FACsPbI<sub>3</sub>. Kinetic traces at a probing wavelength of ~770 nm for e FAPbI<sub>3</sub> and f FACsPbI<sub>3</sub>

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 $FACsPbI_3$  was significantly beneficial for the application of PeLEDs.

#### 3.4 Device Performance of PeLEDs

We fabricated FAPbI<sub>3</sub>- and FACsPbI<sub>3</sub>-based NIR-PeLEDs with the following device configuration: indium tin oxide (ITO) as an anode, poly(ethylene dioxythiophene)/polystyrene sulfonate (PEDOT:PSS) as hole injection layer, poly(9-,9-dioctylfluoreneyl-2,7-diyl)-alt-(4,4'-(N-(4-butylphenyl)diphenylamine) (TFB) as hole transport layer, evaporated perovskite as emission layer, 1,3,5-tris(1-phenyl-1H-benzimidazol-2-yl)benzene (TPBi) as electron transport layer and LiF/Al as cathode (Fig. 5a). The corresponding energy-level diagram of PeLEDs is illustrated in Fig. 5b. To observe the crystal quality and interface, high-resolution scanning transmission electron microscopy (STEM) was performed on the thin cross-section lamella of the PeLED device. The lamellas of devices were prepared via a focused ion beam (FIB) (Fig. S5). As shown in Fig. 5c, FAPbI<sub>3</sub>- and FACsPbI<sub>3</sub>-based PeLED devices were analyzed using STEM under the highangle annular dark-field (HAADF) mode. Based on the STEM images of the FAPbI<sub>3</sub>- and FACsPbI<sub>3</sub>-based PeLED devices in Fig. S6, the thickness of PEDOT:PSS/TFB hole transport layer was determined to be 40~50 nm. The thickness of FAPbI3 and FACsPbI3 emitting layers was determined to be around 20 nm. The thickness of the TPBi electron transport layer was ~ 40 nm. Figure 5d illustrates the energy-dispersive X-ray spectrometer (EDS) elemental mapping of In, I, Cs, Pb, and Al in the FACsPbI<sub>3</sub>-based PeLEDs, demonstrating a uniform distribution of the elements in the PeLED device. Compared with the EDS mapping of FAPbI<sub>3</sub> device in Fig. S6, the only difference observed in the EDS mapping was the Cs element in the FACsPbI<sub>3</sub> emission layer. The current density-voltage and radiance-voltage curves of PeLEDs are shown in Fig. 5e, f. The current density in the FAPbI<sub>3</sub>-based device was observed to be higher than that of FACsPbI<sub>3</sub>-based device (Fig. 5e). The higher current density of FAPbI<sub>3</sub>-based device may be ascribed to the larger grain size of FAPbI3 compared with FACsPbI3, leading to a smaller number of grain boundaries in FAPbI<sub>3</sub> thin films [43]. The FAPbI<sub>3</sub>- and FACsPbI<sub>3</sub>-based devices had a similar radiance of 2.76 and 2.64 W Sr<sup>-1</sup> m<sup>-2</sup>, respectively. As illustrated in Fig. 5g, the EQE value of the device based on FACsPbI<sub>3</sub> reached a maximum of 10.25%, which was approximately 6.5-fold higher than that of the device based on FAPbI<sub>3</sub> (EQE = 1.58%). To address the advantage of the high-throughput optimization of fixing the substrate



Fig. 4 AFM micrographs of a FAPbI<sub>3</sub> and b FACsPbI<sub>3</sub>. FE-SEM images of c FAPbI<sub>3</sub> and d FACsPbI<sub>3</sub>. The crystal size distribution of e FAPbI<sub>3</sub> and f FACsPbI<sub>3</sub>



**Fig. 5** a Device structure, and **b** energy levels of PeLEDs. **c** A STEM-HAADF image of cross-section a full PeLEDs based on FACsPbI<sub>3</sub>. **d** EDS element mapping of Al, Pb, Cs, I, and In of the PeLEDs. Electroluminescence performance: **e** current density–voltage curves, **f** radiance–voltage curves, **g** EQE–current density of FAPbI<sub>3</sub> and FACsPbI<sub>3</sub>. **h** EL spectra of FACsPbI<sub>3</sub>-based PeLEDs at different operating voltages. The inset shows a working device at a voltage of 4 V. Electroluminescence mapping of **i** FAPbI<sub>3</sub>- and **j** FACsPbI<sub>3</sub>-based PeLEDs device at a voltage of 6 V. The yellow dash region and red dash region denoted dark and emission region, respectively

during the evaporation, we also conclude EQE value and wavelength peak distribution of the devices in the different deposited regions. As shown in Fig. S7, the FACsPbI<sub>3</sub>-based devices showed overall enhancements in the efficiency compared with their FAPbI<sub>3</sub>-based counterparts. The efficiency enhancement could be attributed to the enhanced  $\alpha$ -phase crystallinity and reduced trap density of FACsPbI<sub>3</sub>. To provide a comparison, we summarized the performance of stateof-the-art NIR-PeLEDs based on both solution-based process and thermal evaporation (Table S3). The optimal device based on evaporated  $\alpha$ -phase FACsPbI<sub>3</sub> represents one of the best devices for a thermally evaporated NIR-PeLEDs ever reported so far. In order to further investigate the carrier modulation effect, space-charge-limited current (SCLC) analysis was performed. The hole-only devices with a configuration of ITO/PEDOT:PSS/TFB/Perovskite/MoO<sub>3</sub>/Ag and the electron-only devices with a structure of ITO/TPBi/Perovskite/TPBi/LiF/Al based on FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> were fabricated. The trap density ( $n_t$ ) was calculated from the current density–voltage curves using the Mott–Gurney equation [44, 45]:

$$n_t = 2\varepsilon_0 \varepsilon V_{\rm TFL} / eL^2 \tag{3}$$

where  $\epsilon_0$  and  $\epsilon$  are the vacuum permittivity and relative permittivity, respectively. Here, *e* represents the elementary charge and *L* is the thickness of perovskite films. The V<sub>TFL</sub> is the trap-filling limit voltage, which can be estimated by the current density-voltage curves. As shown in Fig. S8, the fitted V<sub>TEL</sub> values of the hole-only device based on FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> were 0.95 and 0.83 V, respectively. Similarly, the V<sub>TFI</sub> values of the electron-only device based on FAPbI<sub>3</sub> and FACsPbI<sub>3</sub> were 1.02 and 0.94 V, respectively. Therefore, the corresponding  $n_t$  values of the hole-only device were estimated to be  $5.25 \times 10^{18}$  and  $4.59 \times 10^{18}$  cm<sup>-3</sup>, respectively. Meanwhile, the  $n_t$  of the electron-only device was calculated to be  $5.64 \times 10^{18}$  and  $5.19 \times 10^{18}$  cm<sup>-3</sup>, respectively. The decreased  $n_t$  value strongly verified the reduced trap densities of the FACsPbI<sub>3</sub> films. As shown in Fig. 5h, the electroluminescence (EL) peak of the device based on FACsPbI<sub>3</sub> was located at ~770 nm. There was no obvious peak shift with the increasing driving voltages. EL mapping was conducted to assess the EL degradation of FAPbI3- and FACsPbI<sub>3</sub>-based PeLEDs. As shown in Fig. 5I, j, a confocal fluorescence microscope was used to measure the EL mapping of the devices at a voltage of 6 V. Due to the photoinactive  $\delta$ -phase and relatively poor morphologies of FAPbI<sub>3</sub>, there was a greatly enlarged dark area (yellow dash region) in the FAPbI<sub>3</sub>-based device during working, which was likely due to the deterioration of the emission layer [46, 47] (Fig. 5i). On the contrary, the FACsPbI<sub>3</sub>-based device exhibited much higher overall radiative intensities and more homogeneous radiative distributions (Fig. 5j).

#### **4** Conclusions

In conclusion, we comprehensively studied the effect of A-site Cs substitution on the optoelectronic properties of FAPbI<sub>3</sub> thin films in vapor deposition. Through triple-source co-evaporation,  $\alpha$ -phase FACsPbI<sub>3</sub> thin films with PL located at 782 nm and a nanocrystal size of~61 nm were prepared. Detailed and systematic characterizations confirmed that partially replacing FA with Cs in the FAPbI<sub>3</sub> can help with the suppression of the non-radiative recombination, reduction of the nanocrystal size, elimination of the metallic Pb, and enhancement of the surface morphology of FAPbI<sub>3</sub> thin films. As a result, the NIR-PeLEDs based on α-phase FACsPbI<sub>3</sub> achieved a maximum EQE of 10.25%, significantly surpassing that of NIR-PeLEDs based on FAPbI<sub>3</sub> (EQE = 1.58%). This work offers a feasible and reliable route to improve the performance of NIR-PeLEDs. Furthermore, it can be expected that this work will also inspire the development of scalable perovskite-based optoelectronic devices.

Acknowledgements This work was financially supported by the "Pioneer" and "Leading Goose" R&D Program of Zhejiang (2024C01192), National Natural Science Foundation of China (62322505, 62425502, 62375276, 62375060, U23A6002, 62050039, 62004075, and 62374069), Shanghai Pilot Program for Basic Research (22JC1403200), and the Natural Science Foundation of Hubei Province (2024AFB423).

Author Contributions Juan Du, Jiajun Luo, and Siwei He conceived the idea of the study, experiment study, and discussion. Siwei He and Lanxin Qin carried out the device fabrications. Siwei He, Qinlanxin, and Zhengzheng Liu carried out the structure characterizations and data analysis. Zhengzheng Liu and Jae-Wook Kang contributed to the formal analysis. Siwei He wrote the manuscript. All authors have given approval to the final version of the manuscript.

#### Declarations

**Conflict of interest** The authors declare no interest conflict. They have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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**Supplementary Information** The online version contains supplementary material available at https://doi.org/10.1007/s40820-025-01776-3.

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