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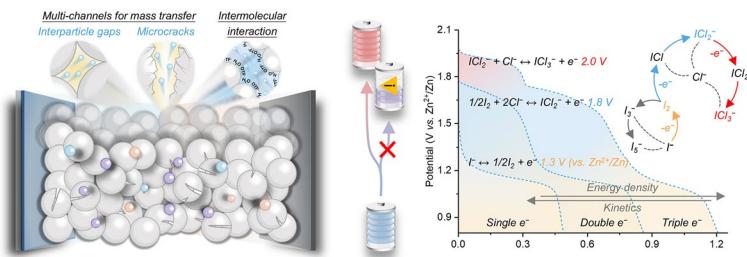
## Electrically Insulating Rigid Multi-Channel Electrolyte Container for Customizable Electron Transfer in Zn-Halogen Batteries

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### HIGHLIGHTS

- A rigid electrolyte container ( $\text{SiO}_2@PVDF-hfp$ ) was designed to immobilize the liquid halogen-ion electrolyte, enabling separator-free Zn-halogen batteries.
- The container regulates  $\text{Zn}^{2+}$  solvation *via* hydrogen bonding regulation, while providing multi-channel structure for enhanced mass transfer, jointly enabling durable Zn plating/stripping.
- Effective confinement of intermediates ensures high reversibility across multi-electron transfer mechanisms, achieving an exceptionally low-capacity decay of 0.02% over 4500 cycles and a high areal capacity of  $11.9 \text{ mAh cm}^{-2}$ .

**ABSTRACT** Recent advancements in Zn-halogen batteries have focused on enhancing the adsorptive or catalytic capability of host materials and stabilizing complex intermediates with electrolyte additives, while the halogen-ion electrolyte modifications exhibit strong potential for integrated interfacial regulation. Herein, we design an electrically insulating rigid electrolyte container to immobilize a liquid halogen-ion electrolyte for separator-free Zn-halogen batteries with customizable electron transfer. Robust hydrogen bonding of hydroxyl groups in  $\text{SiO}_2$  with fluorinated moieties in  $\text{PVDF-hfp}$  regulates  $\text{Zn}^{2+}$  solvation and suppresses  $\text{H}_2\text{O}$  activity, while multi-channels formed by microcracks and interparticle gaps not only enhance mass transfer but also buffer interfacial electric field, jointly enabling a durable Zn plating/stripping. Effective confinement of intermediates also ensures the high reversibility across single- $(\text{I}^-/\text{I}^0)$ , double- $(\text{I}^-/\text{I}^0/\text{I}^+)$ , and triple- $(\text{I}^-/\text{I}^0/\text{I}^+, \text{Cl}^-/\text{Cl}^0)$  electron transfer mechanisms at cathode, as evidenced by the double-electron transfer systems exhibiting a low capacity decay rate



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of 0.02% over 4500 cycles at  $10 \text{ mA cm}^{-2}$  and a high areal capacity of  $11.9 \text{ mAh cm}^{-2}$  at  $2 \text{ mA cm}^{-2}$ . This work presents a novel “container engineering” approach to halogen-ion electrolyte design and provides fundamental insights into the relationships between redox reversibility and reaction kinetics.

**KEYWORDS** Electrolyte container; Interfacial reaction regulation; Customizable electron transfer; Zn-halogen batteries

## 1 Introduction

The past decade has witnessed the explosive growth in Zn-based batteries and their potential for grid-scale energy storage [1]. Distinct from rocking-chair systems driven solely by  $\text{Zn}^{2+}$  transport, the Zn-halogen batteries operate on a dual-ion mechanisms involving independent interfacial reactions (i.e., Zn plating/stripping at anode and halogen redox at cathode) [2]. This configuration offers fast kinetics, high-energy efficiency, and design flexibility, yet the uncontrolled loss of active halogen species leads to short lifespan, thereby limiting their practical applications [3].

Halogens are incorporated into Zn-based batteries primarily through two strategies, including cathode blending and electrolyte additives [4]. The former is more commonly employed for higher utilization of charge carriers, while the latter streamlines the process by focusing optimization on the electrolyte itself [5]. Regarding physicochemical properties, halogen species such as  $\text{I}^-$  and  $\text{Cl}^-$  can form highly polar and unstable interhalogen compounds [6]. For instance, in a common double-electron system,  $\text{I}^-$  can complex with  $\text{I}_2$  to form  $\text{I}_3^-$ , which impairs kinetics. Meanwhile, the oxidized  $\text{ICl}$  species is prone to decomposition, causing the collapse of high-voltage plateau [7]. Hence, designing a system that can effectively confine these reactive species to enable reversible, high utilization redox remains a central goal.

To address these issues, various modification strategies have been proposed. A wide range of carbon substrates, doped frameworks, and metal compound composites were introduced into the cathode to enhance halogen confinement through physical adsorption or chemical anchoring, while catalytic components such as heteroatom dopants and transition metal sites were incorporated to accelerate redox kinetics [8–10]. Additionally, surfactants, organic chelators and pH-responsive polymers were added into electrolyte to stabilize reactive interhalogens via active association or

pH regulation, so as to limit their migration [11–13]. These additives also modulated  $\text{Zn}^{2+}$  solvation and reduced  $\text{H}_2\text{O}$  activity, contributing to inhibited gas evolution and uniform Zn deposition [14–17]. Also, functional separator coatings with ion-selective or adsorption capabilities, such as cation exchange polymers, metal–organic sieves, and Janus structures, were developed to restrict halogen anion migration and prevent species crossover [18]. From another perspective, while strategies often focus on regulating liquid electrolytes within passive glass fiber membranes, a promising alternative is to view the electrolyte and membrane as an integrated functional system. Holistically engineering this entity with synergistic interfaces, chemical constraints and tailored mechanical properties will create a new pathway toward developing high-performance Zn-halogen batteries suitable for diverse applications.

In this work, we move beyond conventional passive filled membranes or *in situ* polymer hydrogels to introduce a new “container engineering” strategy. We fabricate a pre-formed, rigid, and electrically insulating container composed of  $\text{SiO}_2$  and PVDF-*hfp* (denoted as SP), engineered not merely to host the liquid electrolyte (LE) but to actively regulate the electrochemical environment through integrated functions. The obtained rigid film (SPE) enables separator-free Zn-halogen batteries with improved reversibility and lifespan under customizable electron transfer mechanisms. The microcracks and interparticle gaps in SP form multi-channels for ion dissociation, storage, and migration, endowing SPE with ideal mass transfer capability. Meanwhile, strong hydrogen bonding between hydroxyl groups in  $\text{SiO}_2$  and fluorinated moieties in PVDF-*hfp* disrupts the original  $\text{H}_2\text{O}$  network in LE, thus driving out  $\text{H}_2\text{O}$  molecules from  $\text{Zn}^{2+}$  solvation and suppressing their activity. Coupled with buffering effect of multi-channels on ion concentration gradient and electric field, SPE promotes more uniform Zn plating and extended lifespan at a low density of  $0.5 \text{ mA cm}^{-2}$ . Importantly, effective confinement of intermediates localizes redox reactions at cathode interface, achieving high reversibility across the

single- $(I^-/I_0)$ , double- $(I^-/I_0/I^+)$ , and triple- $(I^-/I_0/I^+, Cl^-/Cl_0)$  electron transfer mechanisms. This work presents a new strategy for electrolyte design in Zn-halogen batteries from the perspective of “container engineering” and provides fundamental insights into the correlation between redox reversibility and interfacial reaction kinetics.

## 2 Experimental Section

### 2.1 Materials

All reagents were purchased from Aladdin unless otherwise specified.

### 2.2 Preparations of SP, SPE, and Zn-I Cells

#### 2.2.1 Preparation of $SiO_2@PVDF-hfp$ Disk (Marked as SP)

Firstly, fumed silica ( $SiO_2$ ) and poly(vinylidene fluoride-co-hexafluoropropylene) ( $PVDF-hfp$ ) were mixed in a 4:1 weight ratio in N-methyl pyrrolidone (NMP) and stirred vigorously overnight. The mixture was dried at 60 °C to yield the initial  $SiO_2@PVDF-hfp$  powder. Then, 100 mg of this powder was pressed at 10 MPa for 5 min using a stainless steel mold to form a  $SiO_2@PVDF-hfp$  disk with a diameter of 16 mm. Finally, these disks were further vacuum-dried at 100 °C for 12 h to remove any residual solvent.

#### 2.2.2 Preparation of Aqueous Electrolyte and Electrolyte-absorbing $SiO_2@PVDF-hfp$ Disk (Marked as SPE)

The aqueous electrolyte was prepared by adding 3 M  $Zn(CF_3SO_3)_2$ , 0.5 M potassium iodide (KI), and 2 M potassium chloride (KCl) to water and stirring until the solution became completely clear, among which KI serves as the sole iodine source in this research. The aqueous electrolyte, in its liquid form, is denoted as LE for subsequent comparisons. After immersing the dried SP disk in the aqueous electrolyte for overnight, a transparent SPE disk was obtained, which had absorbed approximately 70  $\mu$ L of LE according to the calculation based on thermogravimetric results (Fig. S6).

#### 2.2.3 Preparation of Electrodes and Fabrication of Zn-I Cells

Zn foil (50  $\mu$ m) without any further modification was used as anode. For the cathode, a slurry was prepared by mixing activated carbon (YP-50F, Kuraray Co., Ltd), Ketjen Black, and polyvinylidene fluoride (PVDF) at an 8:1:1 weight ratio in NMP. The slurry was then uniformly cast onto the current collectors. The approximate activated carbon loading was  $\sim 2$  mg  $cm^{-2}$  when using hydrophilic carbon fiber cloth ( $0.6 \times 0.6$   $cm^2$  in coin cell, W0S1011, Ce Tech Co., Ltd.) and  $\sim 10$  mg  $cm^{-2}$  for cathodes prepared with carbon felt.

CR2025 coin cells assembled in an open air at room temperature were used as battery system. The LE with soaking amount of 100  $\mu$ L was loaded onto a glass fiber separator (GF/D, Whatman), whereas the SPE was used directly without any separator.

In addition, various cells were employed for characterization under different conditions, *e.g.*, Zn/Zn, Zn/Cu, stainless steel/stainless steel cells, pouch cell...

### 2.3 Material Characterizations

X-ray diffractometer (Rigaku Mini Flex 600, Cu  $K\alpha$  radiation,  $\lambda = 1.5418$   $\text{\AA}$ ) was used to determine the crystalline structure and phase composition. Morphology was characterized by a scanning electron microscope (SEM, TESCAN Mira 3), transmission electron microscope (TEM), and high-resolution TEM (HRTEM, FEI Titan G2, 60–300 kV) equipped with an energy-dispersive spectrometer (EDS). Elemental valence states were detected by an X-ray photoelectron spectroscopy (XPS, VG Escalab-250xi). Raman spectra were investigated using Lab RAM HR800, Fourier transform infrared (FT-IR) and UV-vis adsorption spectra were measured by Thermo Fisher Nicolet 6700 and Evolution 220.  $N_2$  adsorption/desorption isotherms and pore size distributions were determined via a Brunauer–Emmett–Teller analyzer (BET, Micromeritics ASAP 2020). Thermogravimetric (TG) analysis was measured by NETZSCH STA 449C. Mercury intrusion porosimetry (MIP, Micro Active Auto Pore V 9600) was carried out to evaluate the pore size distributions ranging from 780  $\mu$ m to 5 nm. Hardness and elastic modulus were measured on Bruker Hysitron TI 980. Differential electrochemical mass spectrometry was collected on Hiden HPR-20 OEMS.

## 2.4 Electrochemical Measurements

All electrochemical performances were conducted on at least five independent cells for each condition to ensure reproducibility. The galvanostatic charge/discharge and Zn plating/stripping measurements were evaluated on a multi-channel battery test system (LAND CT2001A). Cyclic voltammetry (CV), chronoamperometry (CA), Tafel, linear sweep voltammetry (LSV), and electrochemical impedance spectroscopy (EIS) from 100 kHz to 10 mHz were conducted by an electrochemical workstation (CHI660E).

## 2.5 Molecular Dynamics Simulations

Molecular dynamics (MD) simulations were conducted using the Force module in Materials Studio 2023 (Accelrys Inc.). Two models were constructed for comparison, with detailed molecular and ionic compositions summarized in Table 1. All simulations employed the COMPASS III force field. A time step of 1 fs was used. Each system was first equilibrated under the isothermal–isobaric (NPT) ensemble at 300 K and 0.1 GPa by a Berendsen barostat for 2 ns, followed by a 5 ns production run under the canonical (NVT) ensemble using a Nose thermostat. The simulation time was long enough to reach the equilibrium states of the systems. Then, radial distribution functions (RDFs) and hydrogen bonding information were calculated using the VMD analysis tool.

## 2.6 Density Functional Theory Calculations

Density functional theory (DFT) calculations were carried out using the Vienna Ab initio Simulation Package (VASP) [19]. The exchange–correlation interactions were treated with the generalized gradient approximation (GGA) using

the Perdew–Burke–Ernzerhof (PBE) functional, chosen for its reliable balance of computational efficiency and accuracy for quasi-solid-state systems, in conjunction with the projector augmented wave (PAW) method [20]. A vacuum layer of 20.0 Å was applied along the *z*-axis to eliminate the interactions between periodic images. Four-layer slab models were constructed to simulate the adsorption surface, which is a simplified model intended to isolate the intrinsic thermodynamics of the redox reactions, with the bottom two layers fixed to preserve bulk characteristics. A plane-wave cutoff energy of 520 eV was used, and a  $\Gamma$ -centered  $3 \times 1 \times 1$  k-point mesh was adopted. The electronic self-consistent field (SCF) convergence threshold was set to  $1 \times 10^{-5}$  eV atom $^{-1}$  during geometry optimization and tightened to  $1 \times 10^{-7}$  eV atom $^{-1}$  in single-point energy calculations. Then, structural relaxations continued until the residual forces on all atoms were less than 0.02 eV Å $^{-1}$ . Van der Waals interactions were accounted for using the DFT-D4 correction scheme, a modern and accurate method for capturing the interactions critical to adsorption. Then, the Gibbs free energy ( $\Delta G$ ) associated with halogen redox was calculated using the following equation:

$$\Delta G = E_{gs} + E_{zp} - T\Delta S \quad (1)$$

where  $E_{gs}$  is the ground-state energy,  $E_{zp}$  is the zero-point energy, and  $T\Delta S$  represents the entropy term. Both  $E_{zp}$  and  $\Delta S$  were derived from vibrational frequency analyses based on DFT calculations. This overall approach is a standard practice aimed at elucidating qualitative reaction trends rather than replicating the full, complex solvated environment.

## 2.7 COMSOL Simulations

The electric potential and current density distribution during Zn plating were simulated by the Tertiary Current

**Table 1** Molecule numbers in MD simulations (The unit component of PVDF-*hfp* is set as  $(-\text{CH}_2\text{CF}_2-)_4[-\text{CF}_2\text{CF}(\text{CF}_3)-]$ )

Electrolyte	Component					
	$\text{Zn}(\text{CF}_3\text{SO}_3)_2$	KCl	KI	$\text{H}_2\text{O}$	$\text{SiO}_2$	PVDF- <i>hfp</i>
LE	618	412	103	6491	–	–
SPE	264	176	44	1846	930	108

Distribution in COMSOL Multiphysics. The electrochemical reaction kinetics at the electrode surface was described by the Butler–Volmer equation. Crucially, to simulate morphological evolution of the plated layer, Deforming Geometry interface was employed, where the boundary velocity of the electrode surface was set to be proportional to the local current density. The initial geometry included a sinusoidal profile on the electrode surface to represent initial protrusions. The simulations were run for 1000 s to capture the long-term plating behavior, and the corresponding dendrite growth height was extracted for quantitative analysis.

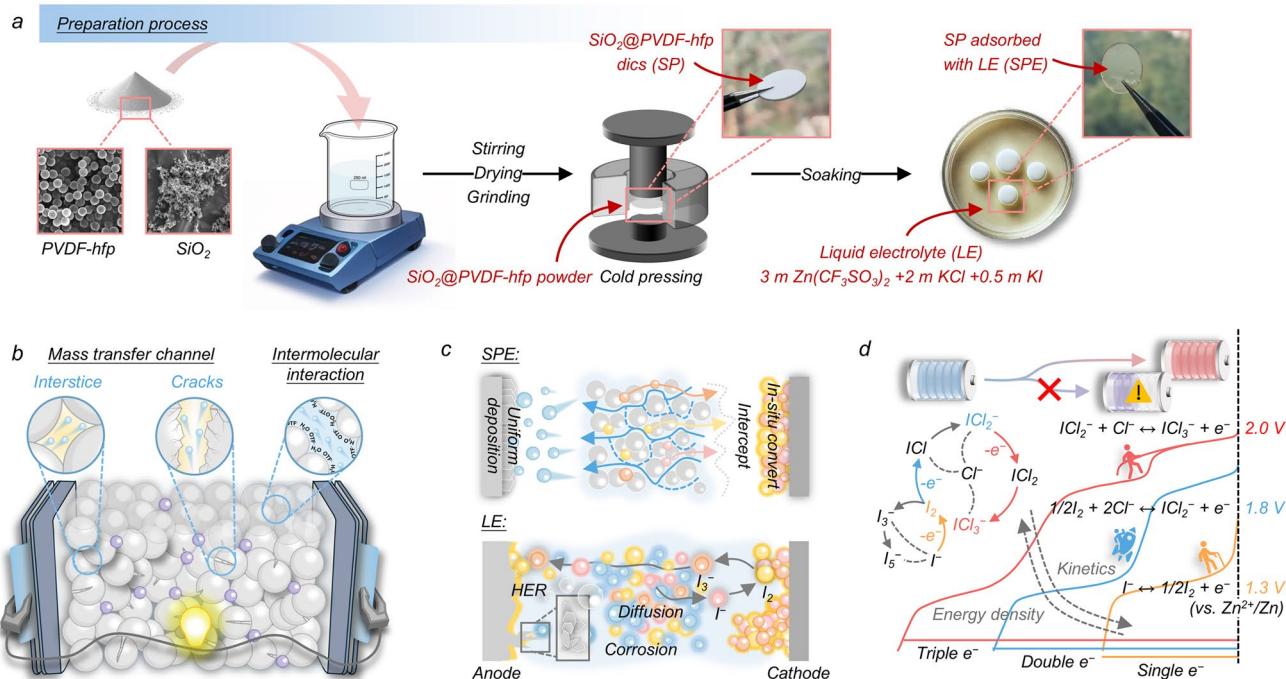
### 3 Results and Discussion

The electrolyte was prepared by fabricating rigid containers followed by a soaking process (Fig. 1a). First, fumed silica and poly(vinylidene fluoride-co-hexafluoropropylene) (PVDF-*hfp*) were vigorously mixed in an organic solvent and dried to form  $\text{SiO}_2$ @PVDF-*hfp* powder, which was then pressed into disk shapes using a mold and further dried to produce mechanically robust  $\text{SiO}_2$ @PVDF-*hfp* (denoted as SP, serving as the rigid container). These disks were

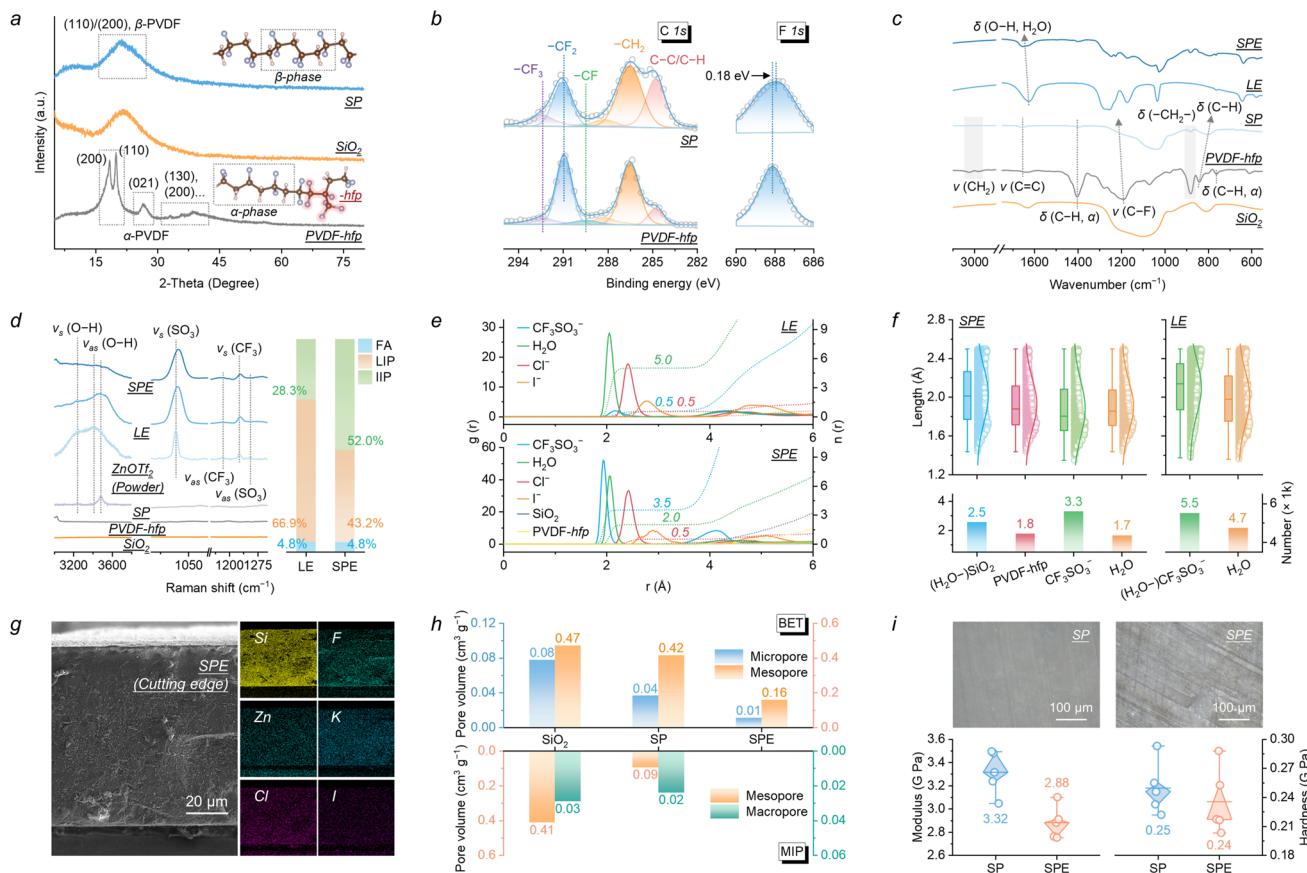
subsequently immersed in a liquid electrolyte (denoted as LE for comparison) for soaking overnight, resulting in transparent electrolyte-adsorbed disks (denoted as SPE, representing the SP-host with electrolyte composite). The presence of potentially interconnected channels, interstices, and cracks in the internal architecture, along with possible interactions between PVDF-*hfp* chains and ions, enhances the ability of such a rigid container to provide storage sites and facilitate mass transfer pathways for the absorbed LE (Fig. 1b). Also, the incorporated framework within SPE helps expel water molecules and suppress their activity, offering advantages in controlling interfacial reactions such as gas evolution at the anode surface and shuttle effect of polyiodide species from the cathode surface (Fig. 1c). This approach provides a novel strategy for designing highly reversible and customized electron transfer Zn-halogen batteries with balanced energy density and kinetics (Fig. 1d).

#### 3.1 Physicochemical Evolution of Preparation

The rigid containers (SP) were prepared by mixing  $\text{SiO}_2$  and PVDF-*hfp* in an N-methyl pyrrolidone (NMP) solvent, followed by drying and pressing, of which the XRD



**Fig. 1** Schematic illustrations of SPE fabrication and functional mechanisms: **a** SPE preparation; **b** mass transport enhancement; **c** interfacial reaction regulation; and **d** customizable charge transfer Zn-halogen battery application



**Fig. 2** Physicochemical characterizations during the preparation: **a** XRD patterns of  $\text{SiO}_2$ , PVDF-*hfp*, and SP; **b** XPS spectra of SP and PVDF-*hfp*; **c** FT-IR spectra of  $\text{SiO}_2$ , PVDF-*hfp*, SP, LE, and SPE; **d** Raman spectra of  $\text{SiO}_2$ , PVDF-*hfp*, SP,  $\text{Zn}(\text{CF}_3\text{SO}_3)_2$ , LE, and SPE; **e** RDFs and H-bonds statistics of LE and SPE derived from MD-simulated solvation boxes; **g** SEM images with EDS elemental mappings of SPE; **h** pore volume distribution of  $\text{SiO}_2$ , SP, and SPE measured via BET and MIP; and **i** average hardness and elastic modulus of SP and SPE calculated from nanoindentation

patterns are shown in Fig. 2a. The sharp diffraction peak of PVDF-*hfp* indicates its high crystallinity, with a doublet at  $18.4^\circ$ ,  $20.0^\circ$ , and a peak at  $26.6^\circ$  corresponding to (020), (110), and (021) crystal planes of the monoclinic  $\alpha$ -phase. Also, weak reflections at  $33.2^\circ$ ,  $35.9^\circ$ ,  $38.8^\circ$ , and  $41.1^\circ$  can be attributed to the (130), (200), (002), and (111) planes, respectively. In contrast, both fumed  $\text{SiO}_2$  and SP display only a broad peak around  $20^\circ$ , indicative of a highly amorphous structure. According to the crystal form of PVDF-*hfp*, this broad peak points to the (110), (200) planes of the orthorhombic  $\beta$ -phase, suggesting a phase transition during this process. In general, the  $\beta$ -phase with TTTT conformation in PVDF-*hfp* exhibits stronger polarity than the  $\alpha$ -phase with TGTG' conformation, which is thought to enhance hydrogen-bonds (H-bonds) formation and ion dissociation (Fig. S1) [21, 22].  $\text{SiO}_2$  exhibits hydrophilicity

due to abundant hydroxyl groups on its surface, which likely form H-bonds ( $-\text{OH}\cdots\text{FC}-$ ) with PVDF-*hfp*. In this scenario, electron transfer (from  $\text{F}^-$  to  $\text{H}^+$ ) alters the binding energy between elements, as detected by X-ray photoelectron spectroscopy (XPS) [23]. The C 1s region after deconvolution by Gaussian fitting reveals the bonding information for  $-\text{CF}_3$ ,  $-\text{CF}_2$ ,  $-\text{CF}$ ,  $-\text{CO}$ ,  $-\text{CH}_2$ , and C–H/C–C groups (Fig. 2b). Compared to PVDF-*hfp*, the peaks for  $-\text{CF}_3$ ,  $-\text{CF}_2$ , and  $-\text{CF}$  in SP show a slight blue shift, together with C–F peak shifts by  $0.18$  eV toward the red in F 1s region, indicating the formation of H-bonds between the hydroxyl groups of  $\text{SiO}_2$  and fluorine groups of PVDF-*hfp* [24]. This is further supported by the enhanced intensity of Si–OH peak in Si 2p region, while the peak for  $\text{Si}^{4+}$  remains unchanged, indicating that  $\text{SiO}_2$  retains its original chemical state in SP (Fig. S2).

After immersing SP in LE overnight and drying, the SPE was obtained. The formation mechanism of this process was further investigated by Fourier transform infrared spectroscopy (FT-IR) and Raman spectroscopy.  $\text{SiO}_2$  exhibits characteristic peaks in FT-IR spectra, such as adsorption peaks at 807 and 1100  $\text{cm}^{-1}$  corresponding to the symmetric stretching ( $\nu_s$ ) of Si–O and antisymmetric stretching ( $\nu_{as}$ ) of Si–O–Si, respectively. The peaks at 970 and 3440  $\text{cm}^{-1}$  for binding vibration ( $\delta$ ) of Si–OH indicate the presence of hydroxyl groups again. Compared to PVDF-*hfp* and  $\text{SiO}_2$ , the adsorption peaks decrease obviously in SP, in which the weakening of  $\nu(\text{CH}_2)$  at 2980  $\text{cm}^{-1}$  demonstrates the  $\text{CH}_2$  deprotonation in PVDF-*hfp*, while the appearance of  $\nu(\text{C}=\text{C})$  at 1660  $\text{cm}^{-1}$  and the blue shift of  $\nu(\text{C–F})$  suggest the dehydrofluorination with mass reduction in PVDF-*hfp* chain [25]. Meanwhile, the red shift of  $\delta(\text{C–H})$  at 840  $\text{cm}^{-1}$  and the almost vanished  $\delta(\text{C–H}, \alpha\text{-phase})$  at 761 and 1403  $\text{cm}^{-1}$  further prove the phase transition from  $\alpha$ -phase to  $\beta$ -phase in the highly polar NMP solvent (Fig. 2c). Immediately after, compared to LE, SPE exhibits a reduced intensity of  $\nu(\text{O–H})$  at 2950–3750  $\text{cm}^{-1}$ , a blue shift of  $\delta(\text{O–H}, \text{H}_2\text{O})$ , and a noticeable narrowing of  $\text{Zn}^{2+}$  solvation-related peak at 752–780  $\text{cm}^{-1}$ , reflecting that the introduction of SP container disrupts the strong H-bonds network in aqueous solution while enhancing the solvation of  $\text{Zn}^{2+}$  with anions (Fig. S3). The ion coordination was further analyzed by Raman spectra, as shown in Figs. 2d and S4. The  $\nu_s(\text{O–H})$  and  $\nu_{as}(\text{O–H})$  peaks at 3240 and 3408  $\text{cm}^{-1}$  of LE disappear in SPE, leaving only a broad hump at 3483  $\text{cm}^{-1}$  associated with structural water, which reveals the transition of water molecules from long-range order to an isolated state [26]. Additionally, the  $\nu(\text{CF}_3)$  and  $\nu(\text{SO}_3)$  at 764 and 1035  $\text{cm}^{-1}$  in  $\text{Zn}(\text{CF}_3\text{SO}_3)_2$  shift to higher frequencies as the material transitions from powder to LE and then to SPE, suggesting enhanced coordination of  $\text{CF}_3\text{SO}_3^-$ . By fitting the spectra in the 750–780  $\text{cm}^{-1}$  region,  $\text{Zn}^{2+}$  association states are classified into free anions (FAs, 757  $\text{cm}^{-1}$ ), loose ion pairs (LIP, 764  $\text{cm}^{-1}$ ), and intimate ion pairs (IIPs, 768  $\text{cm}^{-1}$ ) [27]. As results, the FAs fraction remains unchanged (4.8%) in LE and SPE, while the IIPs fraction enlarges from 28.3% to 52.0% in SPE, confirming stronger solvation of  $\text{Zn}^{2+}$  with  $\text{CF}_3\text{SO}_3^-$  in the SP container.

To explore the differences in solvation chemistry, molecular dynamics (MD) simulations were carried out on both LE and SPE. The 3D snapshots of the solvation boxes show a uniform distribution of all components, with

no evident phase separation or molecular aggregation in the SPE composite (Fig. S5) [28]. Analysis of radial distribution function (RDF) reveals that the SP itself does not participate in the  $\text{Zn}^{2+}$  solvation shell. Instead, it enhances the interaction between  $\text{Zn}^{2+}$  and  $\text{CF}_3\text{SO}_3^-$  anions at the expense of displacing water molecules. This is evidenced by a shift in the average solvation cluster from  $[\text{ZnCl}_{0.5}(\text{CF}_3\text{SO}_3)_{0.5}(\text{H}_2\text{O})_5]^+$  in the LE to the more anion-rich  $[\text{ZnCl}_{0.5}(\text{CF}_3\text{SO}_3)_{3.5}(\text{H}_2\text{O})_2]^{2-}$  in the SPE (Fig. 2e) [29]. This restructuring is driven by changes in H-bonds network. In LE, the network is dominated by  $\text{H}_2\text{O}–\text{H}_2\text{O}$  (46%) and  $\text{H}_2\text{O}–\text{CF}_3\text{SO}_3^-$  (54%) interactions. In SPE, these proportions decrease significantly, with the framework itself forming a substantial number of new H-bonds of  $\text{H}_2\text{O}–\text{SiO}_2$  (27%) and  $\text{H}_2\text{O}–\text{PVDF-}hfp$  (19%) (Fig. 2f; Table S1). Thermogravimetric (TG) analysis reveals a complex, two-stage absorption kinetic-related water molecule for the electrolyte (Fig. S6). A rapid initial uptake of loosely bound “physical water” is followed by a much slower formation of tightly bound “structural water,” which is associated with stable ion solvation sheaths. This reflects that  $\text{SiO}_2$  and PVDF-*hfp* in SP framework not only physically confine the water molecules but also actively restructure the H-bonds network over time. Notably, the clear water content fundamentally distinguishes SPE from solvent-free solid polymer electrolytes.

Physical structural evolution was examined using scanning and transmission electron microscopy (SEM, TEM). The morphology of fumed  $\text{SiO}_2$  consists of a loose network formed by fused nanoparticles approximately 10 nm in size, while PVDF-*hfp* exhibits uniform nanospheres with a diameter of ~300 nm. The SP container exhibits a compact and flat surface with a uniform elemental distribution. This structure is achieved as the phase transition in the NMP solution causes the PVDF-*hfp* spheres to collapse and then act as a binder under external pressure, forming a dense network of the  $\text{SiO}_2$  nanoparticles (Fig. S7). Furthermore, there are mesopores and macropores formed between  $\text{SiO}_2$  particles at the SP surface. Interstices ranging from nanometers to micrometers in length and fissures with widths approaching 500 nm were observed at both the edges and internal cross-sections, serving as multi-channels for transporting LE (Fig. S8). The SPE obtained after soaking appears transparent, and its morphology remains unchanged upon a heat treatment. Also, elemental mappings by energy-dispersive X-ray spectroscopy (EDS) of the surface, edge, and internal cross-section reflect the uniformity of ion distribution

after LE penetration (Figs. 2g, S9). Inspired by these, Brunauer–Emmett–Teller (BET) test and mercury intrusion porosimetry (MIP) were carried out, as shown in Fig. 2h. The specific surface area of  $\text{SiO}_2$  decreases from 176.85 to  $86.99 \text{ m}^2 \text{ g}^{-1}$  after being mixed and pressed with PVDF-*hfp* ( $11.12 \text{ m}^2 \text{ g}^{-1}$ ), while the pore volume reduces slightly from 0.47 to  $0.42 \text{ m}^3 \text{ g}^{-1}$ , and the micropore distribution enlarges slightly. This is attributed to the closer contact between the nanoparticles under pressure as well as the formation of new micropores under collapse. Mesopores are formed in the cracks and gaps between the compact particles during this period, making the mesopores dominant in SP. While in the dried SPE, it shows a further decrease in both the specific area and pore volume, with the contents of micropores, macropores, and macropores unchanged significantly. This means that the LE infiltration does not damage the pore structure, and its crystallization can be extensively confined in SP container (Figs. S10, S11). In general, high specific area and narrow space built by mesopores are conducive for ion(s) adsorption and storage in SPE, while the ordered channels of macropores enhance the mass transfer [30, 31]. The nanoindentation test using Oliver–Pharr method was conducted to evaluate the mechanical properties [32]. The SP exhibits a hardness and elastic modulus of 0.25 and 3.32 GPa, respectively, which remained high at 0.24 and 2.88 GPa for SPE (Figs. 2i, S12). This retained high rigidity, coupled with ideal dimensional stability evidenced by a thickness swelling of less than 2%, suggests that SPE processes significant potential to mechanically inhibit Zn dendrites penetration (Fig. S13) [33].

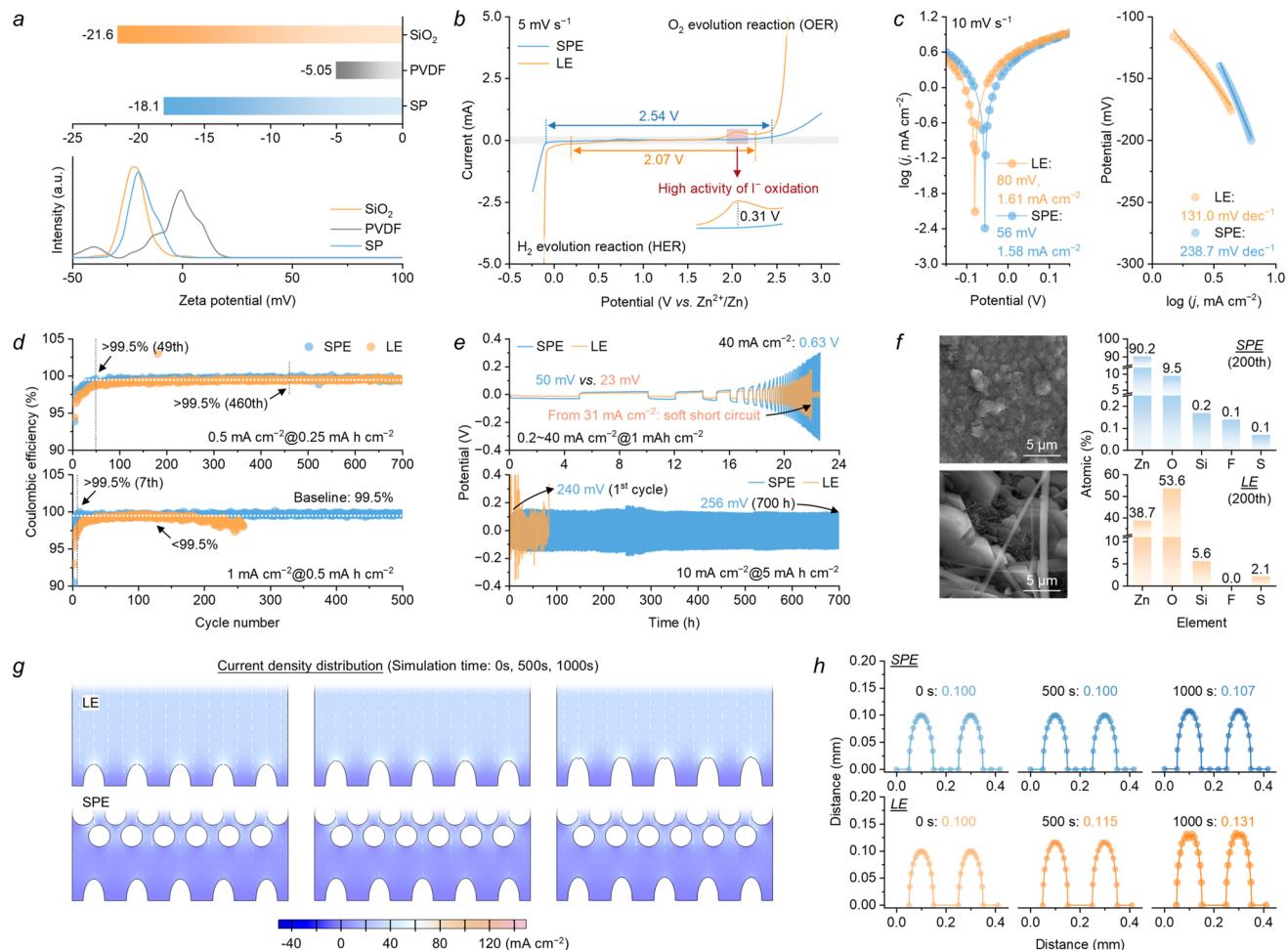
Above all, the mechanically reinforced SPE demonstrates promising mass transfer characteristics. Chemically, the preserved H-bonds network between  $\text{SiO}_2$  hydroxyl groups and PVDF-*hfp* fluorinated moieties in SP container regulate  $\text{Zn}^{2+}$  solvation of introduced LE by disrupting H-bonds networks of water molecules as well as enhancing its interactions with  $\text{CF}_3\text{SO}_3^-$ . Physically, the microcracks and interparticle voids within  $\text{SiO}_2$  matrix establish multi-channels, enabling efficient ion dissociation, temporary storage, and directional migration.

### 3.2 Regulation of Anode Surface Reaction

The zeta potential of SP and each component reflects their internal particle distribution and surface charge

characteristics, where SP exhibits a slightly lower absolute  $\zeta$  potential than  $\text{SiO}_2$  ( $-18.1$  vs.  $-21.6$  mV) while maintaining a high magnitude, indicating the robust electrostatic repulsion and good dispersibility of its particles. Combined with its electronegativity-driven cation capture capability, SP demonstrates potential advantages in optimizing interfacial contact and charge transfer efficiency (Fig. 3a) [34, 35]. This difference is further reflected in dielectric spectra, where SP exhibits a higher dielectric constant and lower dielectric loss than  $\text{SiO}_2$  in the low-frequency range, indicating its enhanced ionic dissociation and subsequent ion transport (Fig. S14) [36]. The  $\text{Zn}^{2+}$  transference number ( $t_{\text{Zn}^{2+}}$ ) and ionic conductivity ( $\sigma$ ) of LE and SPE, measured in Zn/Zn and stainless steel/stainless steel cells for revealing the differences in mass transfer capabilities between liquid- and quasi-solid-state electrolytes, show  $\sigma$  values of 0.058 and  $0.224 \text{ S cm}^{-1}$  and  $t_{\text{Zn}^{2+}}$  values of 0.58 and 0.70 in SPE and LE, respectively (Figs. S15, S16). The diminished ionic conductivity of SPE stems from the electrically insulating  $\text{SiO}_2$ , yet its ion transport performance remains superior to most reported quasi-solid-state electrolytes owing to its internally rich H-bonds network and multi-channels formed between interconnected particles [37, 38]. Linear sweep voltammetry (LSV) at  $5 \text{ mV s}^{-1}$  reveals an extended electrochemical stability window for SPE (2.54 V) versus LE (2.07 V), with SPE showing a flattened slope at both ends, indicating suppressed gas evolution and eliminating redox peaks in the iodine conversion region (Fig. 3b). Such a phenomenon was further confirmed by Tafel curves at  $10 \text{ mV s}^{-1}$ , in which SPE exhibits lower exchange current density ( $1.58$  vs.  $1.61 \text{ mA cm}^{-2}$ ), higher equilibrium potential (56 vs. 80 mV), and larger Tafel slope ( $238.7$  vs.  $131.0 \text{ mV dec}^{-1}$ ) than that in LE (Fig. 3c). This demonstrates that SPE effectively enhances the anti-corrosion capability on anode surface, especially in hydrogen evolution reaction (HER) suppression. This was unequivocally confirmed by *in situ* differential electrochemical mass spectrometry (DEMS), which quantifies a dramatically lower  $\text{H}_2$  evolution rate for SPE system compared to LE in symmetric cells (Fig. S17).

The Zn/Cu cells were assembled to investigate the improvement of Zn plating/stripping behavior, in which the SPE enabled a Coulombic efficiency (CE) of 99.5% within just 49 cycles at  $0.5 \text{ mA cm}^{-2}$ , a milestone that required 460 cycles in LE (Fig. 3d). Such a performance was highly



**Fig. 3** Investigation of anode surface reaction: **a** zeta potentials of  $\text{SiO}_2$ ,  $\text{PVDF-}h\text{fP}$ , and  $\text{SP}$ ; **b** LSV curves at  $5 \text{ mV s}^{-1}$  of  $\text{LE}$  and  $\text{SPE}$ ; **c** Tafel curves at  $10 \text{ mV s}^{-1}$  with Tafel slopes of  $\text{LE}$  and  $\text{SPE}$ ; **d** CE plots at  $0.5 \text{ mA cm}^{-2}$  at  $0.25 \text{ mAh cm}^{-2}$  and  $1 \text{ mA cm}^{-2}$  at  $0.5 \text{ mAh cm}^{-2}$  in  $\text{Zn/Cu}$  cells with  $\text{LE}$  and  $\text{SPE}$ ; **e** current density tolerance between  $0.2$  and  $40 \text{ mA cm}^{-2}$  to  $1 \text{ mA cm}^{-2}$  and lifespan at  $10 \text{ mA cm}^{-2}$  at  $5 \text{ mAh cm}^{-2}$  of  $\text{Zn-Zn}$  cells with  $\text{SPE}$  and  $\text{LE}$ ; **f** SEM images with EDS elemental ratios of deposited  $\text{Zn}$  anode after 200 cycles at  $0.5 \text{ mA cm}^{-2}$  at  $0.25 \text{ mAh cm}^{-2}$  in  $\text{LE}$  and  $\text{SPE}$ ; **g** current density distribution; and **h** dendrite height statistics over time in  $\text{LE}$  and  $\text{SPE}$  by COMSOL simulation

reproducible, as demonstrated by the excellent cell-to-cell consistency shown in Fig. S18. Even at  $1 \text{ mA cm}^{-2}$ , the  $\text{SPE}$  maintained high reversibility, whereas the CE in  $\text{LE}$ -based cell deteriorated. While the  $\text{SPE}$  exhibited higher voltage hysteresis due to its moderate conductivity, this was a worthwhile trade-off for the dramatic improvement in stability and CE (Fig. S19). This advantage was further demonstrated in  $\text{Zn/Zn}$  cells under more demanding conditions. The  $\text{SPE}$  showed exceptional rate tolerance up to  $40 \text{ mA cm}^{-2}$ , in stark contrast to the  $\text{LE}$ -based cell which suffered a “soft short-circuit” above  $31 \text{ mA cm}^{-2}$  (Fig. 3e). Furthermore, it achieves a remarkable lifespan of over 700 h at a demanding

condition of  $10 \text{ mA cm}^{-2}$  to  $5 \text{ mAh cm}^{-2}$ , whereas the cell with  $\text{LE}$  failed in the initial cycles (Fig. S20).

To understand the origin of this enhanced stability, the deposited morphology and composition on the anode were investigated. Post-mortem analyses were conducted on anodes cycled at a low current density of  $0.5 \text{ mA cm}^{-2}$ , a condition where chemically-driven parasitic reactions such as gas escape and by-product formation are particularly prominent [39]. *Ex situ* SEM revealed a dense, flat morphology for the electrode from  $\text{SPE}$ , in contrast with the disordered flake stacking observed from  $\text{LE}$  (Figs. S21, S22). EDS mappings with quantitative analysis showed that the oxygen content on  $\text{SPE}$ -cycled anode was only 9.4%,

far lower than 53.6% on LE-cycled anode (Fig. 3f). This enhanced chemical stability directly translates to a more uniform physical Zn plating, as visualized by *in situ* optical microscopy (Fig. S23). While distinct convex protrusions appeared within 5 min in LE and evolved into irregular “dendritic structures,” the electrode in SPE remained flat throughout the measurement, showing only a slight and uniform thickening. COMSOL simulations provide modeled current density distributions in free- and channel-confined mass transfer (Fig. 3g). Initial surface protrusions trigger a classic “tip effect” in LE system, leading to a severely focused current density with peak value at the tips reaching over 5 times the surrounding valley. This theoretical current focusing is then quantified by the derived dendrite height statistics, which show a rapid, accelerating growth rate characteristic of uncontrolled plating (Fig. 3h). Furthermore, these current “hot spots” are regions of high local overpotential that drive parasitic side reactions like HER, which is in good agreement with our Tafel analysis. *Ex situ* XRD patterns revealed that the electrodes from SPE maintained a pure Zn phase over 200 cycles, while LE counterparts exhibited additional hydroxides ( $Zn_5(OH)_8Cl_2 \cdot H_2O$ ) and oxides (ZnO), due to the release of hydroxide ions from these side reactions (Fig. S24). In stark contrast, the multi-channel structure in SPE acts as an effective ion flux homogenizer, maintaining a remarkably uniform current density, which provides a clear theoretical explanation for the slow, suppressed dendrite growth and ultimately underpins the experimentally observed clean, by-product-free anode surface and the pure Zn phase in the XRD patterns.

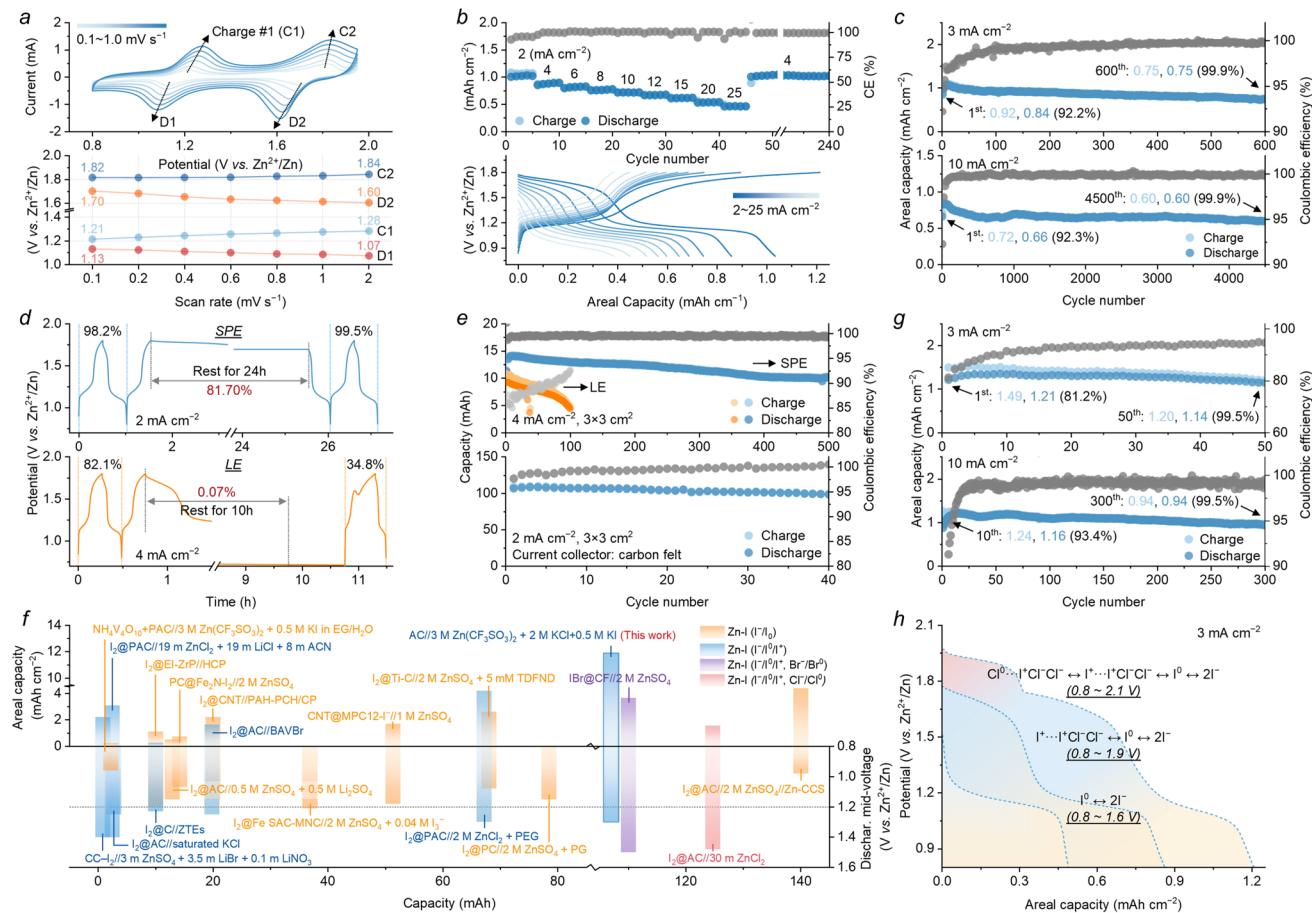
The aforementioned results prove the advantages of SP in regulating the anode surface reaction. The electronegative network formed by dispersed  $SiO_2$  as well as disruption of the original H-bonds network of water in LE efficiently drives  $Zn^{2+}$  transport while suppressed gas evolution reaction, thereby enhancing the reversibility and stability. In addition, SPE with mechanical strength and multi-channels buffers electric fields and concentration gradients between interfaces, promoting uniform Zn plating.

### 3.3 Application of SPE in Customizable Electron Transfer Zn-Halogen Batteries

Due to chloride-induced corrosion of stainless steel under electric fields, most reported Zn-halogen cells with

multi-electron transfer rely on Swagelok-type devices [6, 7]. Since the strong ion adsorption with dissociation and suppressed water activity in SP, SPE holds potential as an electrolyte for multi-electron transfer Zn-halogen coin cells, with potassium chloride and potassium iodide in the soaking solution serving as the only driving forces of this reaction (Fig. S25). Cyclic voltammetry (CV) curves at scan rates of  $0.1\text{--}1.0\text{ mV s}^{-1}$  were firstly collected within the voltage range from 0.80 to 1.95 V, revealing two pairs of redox peaks ( $1.21/1.13\text{ V}$  and  $1.82/1.70\text{ V}$  at  $0.1\text{ mV s}^{-1}$ ) corresponding to a double-electron transfer of  $I^- (I^- \rightarrow I^0 \rightarrow I^+)$  without any accompanying side reactions. Notably, these peaks shifted slightly as the scan rate increased, indicating rapid reaction kinetics at both interfaces (Fig. 4a). Fitting the peak current and scan rate yielded the  $b$ -values around 0.4–0.7 in these redox peaks, revealing a combined diffusion–capacitive-controlled behavior. While quantitative convolution showed capacitance gradually dominating the capacity contribution with increased scan rates (from 40% at  $0.1\text{ mV s}^{-1}$  to 85% at  $1\text{ mV s}^{-1}$ ), further confirming enhanced kinetics and efficient energy storage (Fig. S26).

In terms of the rate performance, SPE delivers average areal capacities of  $1.02\text{ mAh cm}^{-2}$  at  $2.0\text{ mA cm}^{-2}$  and  $0.46\text{ mAh cm}^{-2}$  at  $25\text{ mA cm}^{-2}$ , with Coulombic efficiency (CE) improving from 95.5% to nearly 100% (Fig. 4b). Upon current density recovered to  $4\text{ mA cm}^{-2}$ , the cell operates stably for an additional 200 cycles. Regarding lifespan, capacity decay rates are only 1.8% over 600 cycles at  $3\text{ mA cm}^{-2}$  and 0.02% over 4500 cycles at  $10\text{ mA cm}^{-2}$ , respectively, with a distinct plateau observed at high current density. A capacity fading primarily occurs in the first 100 cycles, during which an equilibrium interface forms concurrently, driving the CE to reach 99.5% (Figs. 4c, S27). Inevitably, the  $I^0/I^+$  redox is inhibited by polarization at  $20\text{ mA cm}^{-2}$ , resulting in a single-electron transfer that inhibits highly reproducible stability, with no capacity decay and CE of nearly 100% over 4000 cycles (Fig. S28). This phenomenon reflects a design trade-off between ionic conductivity and interfacial reaction control. While SPE’s moderate conductivity inevitably increases polarization, its key advantage is enabling stable operation under demanding current densities. Resting tests from the fully charged state were conducted to assess the stability of oxidation products in LE and SPE, in which SPE maintains a CE of 81.7% after standing for 24 h, whereas LE drops sharply to 0.07% within just 10 h (Fig. 4d). Overall, SPE demonstrates superior redox kinetics as well as



**Fig. 4** Electrochemical performances of SPE in customizable charge transfer Zn-halogen batteries: **a** CV curves with redox peak voltages at 0.1–1.0 mV s<sup>-1</sup>; **b** rate capability with voltage curves at 2–25 mA cm<sup>-2</sup>; **c** lifespans at 3 and 10 mA cm<sup>-2</sup> and **d** resting tests (vs. LE) in coin cells with double-electron transfer; **e** lifespans using a carbon fiber cloth current collector at 4 mA cm<sup>-2</sup> (vs. LE) and a carbon felt current collector at 2 mA cm<sup>-2</sup> in pouch cells (3×3 cm<sup>2</sup>) with double-electron transfer; **f** comparisons of areal capacity, discharge mid-voltage, and capacity in other metal-halogen batteries; **g** lifespans at 3 and 10 mA cm<sup>-2</sup> in coin cells with triple-electron transfer; and **h** summarized performances of SPE in customizable charge transfer Zn-halogen coin cells

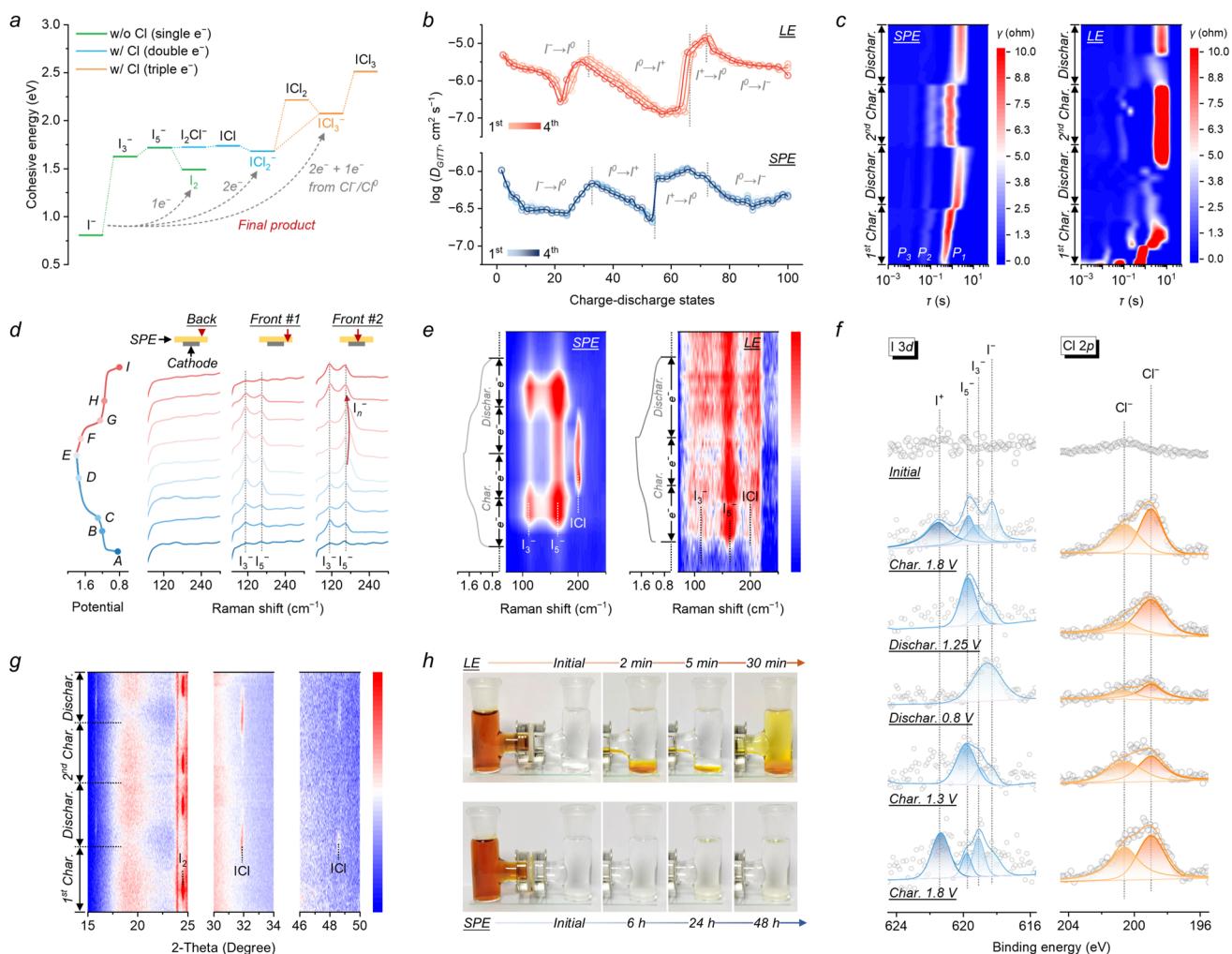
reaction reversibility and stability in coin devices. Furthermore, pouch cells (3×3 cm<sup>2</sup>) were assembled, in which SPE performs a discharge capacity of 10.3 mAh with a retention of 93% over 400 cycles at 4 mA cm<sup>-2</sup>. In contrast, LE shows a lower initial capacity of 9.3 mAh, poor CE (< 90%), and a significantly reduced retention of 46.2% over 100 cycles, indicating limited reversibility of iodine redox in aqueous electrolyte. To further enhance the areal capacity, the carbon fiber cloth current collector was replaced with carbon felt. Owing to its three-dimensional, non-woven architecture, the carbon felt supports a higher loading of activated carbon and provides a vastly expanded reaction interface for enhancing the reaction depth. As a result, a pouch cell (3×3 cm<sup>2</sup>) achieves a significantly improved capacity of 107 mAh, equivalent to an areal capacity of 11.9 mAh cm<sup>-2</sup>,

while remaining stable for over 40 cycles at 2 mA cm<sup>-2</sup> (Figs. 4e, S29). Compared to the reported Zn-halogen pouch cells, SPE demonstrates a clear advantage in areal capacity and exhibits a higher discharge mid-voltage among Zn-I pouch cells with double-electron transfer due to its reduced polarization (Fig. 4f; Table S4).

To evaluate the suitability of SPE in facilitating chloride redox (Cl<sup>-</sup>/Cl<sup>0</sup>), the charging cutoff voltage of coin devices was raised to 2.1 V. While operation at this high voltage raises concerns about stability, the process proceeds via a non-gas-evolving interhalogen pathway (i.e., from ICl<sub>2</sub><sup>-</sup> to ICl<sub>3</sub><sup>-</sup>). This conclusion is clearly supported by in situ DEMS, which detected no Cl<sub>2</sub> signals during high-voltage charging (Fig. S30). Additionally, post-mortem analysis confirms that SPE effectively mitigates

the associated corrosion risk (Fig. S31) [40]. As a result, the cell delivers an initial areal capacity of  $1.21 \text{ mAh cm}^{-2}$  with a lifespan of 50 cycles at  $3 \text{ mA cm}^{-2}$ . Even at  $10 \text{ mA cm}^{-2}$ , a peak capacity of  $1.16 \text{ mAh cm}^{-2}$  with 81% retention was achieved over 300 cycles (Fig. 4g). The robustness of this long-term stability was validated through tests on multiple cells, which exhibited excellent consistency (Fig. S32). A distinct voltage plateau for  $\text{Cl}^-/\text{Cl}^0$  redox observed at both current densities reveals that the excellent kinetics of SPE are maintained in a triple-electron transfer process (Fig. S33). Compared to reported metal-halogen coin cells, SPE-supported cells

exhibit significant improvements in capacity, lifespan, and reversibility across multiple electron transfer mechanisms (Fig. S34). Additionally, a boarder comparison with optimized strategies involving Janus separators, hydrogels, and quasi-solid electrolytes for Zn-based batteries underscores the superior overall balance of properties in our system (Table S5). At this stage, the SPE with effective interfacial reaction control and optimized kinetics characteristics has promising potential as an electrolyte for Zn-halogen cells with customizable charge transfer (Fig. 4h).



**Fig. 5** Reversibility and kinetics of cathode surface reaction: **a** Cohesive energies of interhalogen species involved in multi-electron transfer; **b** diffusion coefficients ( $D_{GITT}$ ) calculated from GITT at  $3 \text{ mA cm}^{-2}$  and **c** DRT spectra at timescales of  $10^{-3}$  to  $50 \text{ s}$  derived from *ex situ* EIS measurements in full cell with LE and SPE; **d** *ex situ* Raman spectra of SPE at different positions in full cell; **e** *in situ* Raman spectra of cathode in full cell with LE and SPE; **f** *ex situ* XPS spectra and **g** *in situ* XRD patterns of cathode in full cell with SPE; and **h** optical images of poly-iodide diffusion in SPE and LE within an H-type device

### 3.4 Reversible and Kinetic Optimizations of Cathode Surface Reaction

Generally, multi-electron transfer in iodine species can be activated and controlled by regulating the stability of interhalogen compounds [7]. Given the diversity of these intercompounds, the pathways of halogen redox in this system were inferred by calculating the potential energies of possible products (Fig. 5a). For single-electron transfer, the complete oxidation product of  $I^-$  is  $I_2$ , which exhibits lower energy compared to other polyiodide ions and limited solubility in aqueous solvent. In terms of double-electron transfer, forming  $ICl_2^-$  is energetically more favorable than  $ICl$ , resulting in a greater tendency for  $I^+$  to coordinate with two  $Cl^-$  [41]. Subsequently,  $ICl_2^-$  is further oxidized into  $ICl_2$  and ultimately stabilized as  $ICl_3^-$ , in which  $Cl^0$  is immobilized through chemical bonding in interhalogen compound rather than physical adsorption [40].

With the halogen redox pathway clarified, kinetic differences and interfacial control between SPE and LE were evaluated based on the double-electron transfer mechanism. Galvanostatic intermittent titration technique (GITT) measurement at  $3\text{ mA cm}^{-2}$  was performed to reveal kinetic characteristics. Overall, the similar trends of  $D_{\text{GITT}}$  across different cycles reflect consistent kinetic behaviors in LE and SPE, with the  $D_{\text{GITT}}$  values in LE being approximately one order of magnitude higher than those in SPE owing to the intrinsic mass transfer advantages of aqueous solvent. Specifically, the  $D_{\text{GITT}}$  in LE exhibits a significant decline in both the  $I^-/I^0$  and  $I^0/I^+$  stages, where the former represents an energy barrier arising from the liquid–solid-phase transition, and the latter results from kinetics deterioration caused by excessive  $ICl$  formation [42]. The spontaneous reduction of highly reactive  $ICl$  also triggers a sharp increase in  $D_{\text{GITT}}$  upon the subsequent discharging. In contrast, SPE displays a similar trend but with notably reduced fluctuation, suggesting that SPE offers a buffering effect on kinetics behavior (Figs. 5b, S35).

To gain deeper insights, *ex situ* electrochemical impedance spectroscopy (EIS) was collected within the first two cycles. The bulk resistance ( $R_b$ ) in SPE shows a slight increase as charging, while charge transfer resistance ( $R_{ct}$ ) remains stable. No additional semicircles point to mass transfer between SPE and cathode, further supporting the fast kinetic behavior. In contrast, LE exhibited initially higher  $R_b$  that gradually decreases with cycling,

accompanied by irregular data points indicative of complicated internal reactions and side products (Fig. S36). Thereafter, the represented distribution relaxation times (DRT) were calculated with the results showing three relaxation times ( $\tau$ ) in the scales of  $10^{-3}$ – $10$  s, which are assigned to ion transport within the electrolyte ( $P_3$ ), charge transfer during redox reactions of iodine species ( $P_2$ ), as well as mass transport of products at the interface and within the cathode ( $P_1$ ) [43–45]. The variations in  $g(\tau)$  suggest that the overall kinetics in both SPE and LE are primarily influenced by  $P_1$  process. Specifically, SPE exhibits shorter timescales in  $P_1$  process, along with consistently lower  $g(\tau)$  in  $P_2$  and  $P_3$  processes during charging, demonstrating accelerated iodine conversion at interface as well as rapid adsorption–desorption by the porous activated carbon, particularly during oxidation. In contrast, LE displays higher initial  $g(\tau)$  across all three processes, in agreement with its elevated  $R_b$  and  $R_{ct}$  in EIS results, which subsequently decline rapidly following redox progression. Notably,  $g(\tau)$  for  $P_1$  process significantly enlarges again during second charging, attributable to residual solid-phase products (e.g.,  $I_2$ ) from incomplete reduction impeding subsequent mass transfer. Irregular timescale distributions of  $P_2$  and  $P_3$  further highlight the complex internal kinetics within LE (Fig. 5c).

Energy storage mechanism and reversibility were confirmed by monitoring iodine species evolution at different interfaces within the battery. *Ex situ* Raman spectra at various positions in SPE were carried out, in which typical scattering peaks of  $I_3^-$  and  $I_5^-$  are located at  $113$  and  $164\text{ cm}^{-1}$ , respectively (Fig. 5d). Throughout cycling, no iodine-related peaks are detected on the back side of SPE (adjacent to anode, marked as *Back*), whereas the front side away from cathode (*Front #1*) displays reversible  $I_3^-$ ,  $I_5^-$  signals with former dominating in intensity. This phenomenon is also observed on the front side aligned with cathode (*Front #2*), but here the  $I_5^-$  signal exhibits higher intensity and a blue shift at high voltage, indicative of higher-order polyiodide ions ( $I_n^-$ ) formation. These results suggest that SP framework effectively regulates the iodine concentration gradient, strictly confining polyiodide ions at the cathode surface, thus preventing interfacial crosstalk and charge carrier loss. Then, *in situ* Raman was used to investigate the differences in cathode between SPE and LE, in which the  $I_3^-$  and  $I_5^-$  signals, along with a  $ICl$  signals at  $200\text{ cm}^{-1}$  sequentially appears and subsequently disappear in SPE, highlighting superior

reversibility. In contrast, these signals persist continuously in LE, reflecting limited controllability (Fig. 5e).

*Ex situ* XPS spectra of cathodes were analyzed to clarify the evolution of chemical valence states. Initially, no peaks appeared in I 3d and Cl 2p regions. Upon charging to 1.8 V, a series of peaks emerge in I 3d region, which can be deconvoluted by Gaussian–Lorentzian fitting into sub-peaks representing I<sup>+</sup>, I<sub>5</sub><sup>−</sup>, I<sub>3</sub><sup>−</sup>, and I<sup>−</sup>, with the integrated area of I<sup>+</sup> dominating. Concurrently, the Cl 2p region shows the emergence and intensification of distinct peaks corresponding to Cl<sup>−</sup> upon charging to 1.8 V. This indicates that Cl<sup>−</sup> does not participate directly in the redox reaction but instead migrates to and accumulates at the cathode interface. This accumulation of Cl<sup>−</sup> is consistent with the formation of I<sup>+</sup>, providing the necessary counter-ions to form interhalogen complexes such as ICl<sub>2</sub><sup>−</sup>, the formation of which is supported by the Gibbs free energy calculations [46]. During discharge, these peaks gradually shift toward lower binding energies, reflecting a reduction in valance states. I<sup>+</sup> peak disappears at 1.25 V, accompanied by increased intensity of I<sub>3</sub><sup>−</sup> and I<sub>5</sub><sup>−</sup>. Ultimately, only I<sup>−</sup> and minor amounts of Cl remain after full discharge. This evolution repeats in subsequent charging, confirming the reversible double-electron transfer of iodine species at the cathode surface (Fig. 5f). Meanwhile, C 1s region remains unchanged throughout, indicating only physical adsorption role of porous activated carbon in this period (Fig. S37). On the other hand, *in situ* XRD patterns were performed to investigate phase evolution on cathode. A characteristic peak at 24.4°, corresponding to (111) crystal plane of I<sub>2</sub>, appears at the initial stage of charge. This peak vanishes at the high-voltage range, accompanied by the emergence of a pair of peaks at 31.9° and 48.6°, attributed to the ICl phase. During subsequent discharge, the ICl phase disappears with the re-emergence of the I<sub>2</sub> phase. Such a reversible evolution was consistently observed over two cycles, confirming the reversibility of iodine redox at the cathode surface (Fig. 5g).

The previously mentioned confinement effect of SP framework on polyiodide ions was validated by resting tests in an H-type cell. As a result, the polyiodides rapidly diffuse through the glass fiber separator within 30 min, whereas SPE setup effectively blocks their migration for at least 48 h (Fig. 5h). Such a difference also pays for distinct Zn plating behaviors on anode surface, as evidenced by SEM images. The anode in SPE exhibits a flat and uniform morphology even after 100 cycles, while fiber residues from the separator and an oxygen

distribution from active water molecules are observed in LE, aligning with earlier conclusions (Fig. S38).

Thus, a clear correlation between reversibility and kinetics of cathode surface reaction is established. SPE enhances reversibility by confining polyiodide species and stabilizing interhalogens, which assist in buffering redox fluctuations and further ensuring stable mass transport. In contrast, uncontrollable and complex internal dynamics in LE lead to poor sustainability, where the accumulation of by-products further impairs reaction kinetics.

#### 4 Conclusions

Overall, an electrically insulating electrolyte container composed of SiO<sub>2</sub> and PVDF-*hfp* (SP) was developed to host liquid halogen-ion electrolyte (LE), forming a rigid film (SPE) that enables separator-free Zn-halogen batteries with reversibly customizable electron transfer. Multi-channels supported by microcracks and interparticle gaps provide pathways for ion dissociation, storage, and transport, while strong hydrogen bonding of hydroxyls in SiO<sub>2</sub> with fluorinated moieties in PVDF-*hfp* disrupts original H<sub>2</sub>O network in LE, thereby modulating the Zn<sup>2+</sup> solvation as well as buffering concentration gradients and electric fields to promote uniform Zn plating. Meanwhile, SPE demonstrates effective confinement of intermediates, achieving high reversibility and extended lifespan in Zn-halogen batteries across the single-(I<sup>−</sup>/I<sup>0</sup>), double-(I<sup>−</sup>/I<sup>0</sup>/I<sup>+</sup>), and triple-(I<sup>−</sup>/I<sup>0</sup>/I<sup>+</sup>, Cl<sup>−</sup>/Cl<sup>0</sup>) electron transfer mechanisms. Although SPE exhibits slightly weaker mass transfer characteristics than LE, it significantly enhances spatial isolation of polyiodides and interhalogen species for maintaining kinetic stability. Whereas aqueous electrolytes like LE suffer from uncontrollable redox dynamics and by-products accumulation, leading to kinetic deterioration and limited lifespan. This work offers a novel approach for halogen-ion electrolyte design from the perspective of container engineering and provides valuable insights into the interplay between redox reversibility and interfacial reaction kinetics.

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#### Declarations

**Conflict of Interest** The authors declare no conflict of interest. They have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper. Prof. Jiang Zhou is an editorial board member for *Nano-Micro Letters* and was not involved in the editorial review or the decision to publish this article.

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