

# MoS<sub>2</sub> Nanosheet Arrays Rooted on Hollow rGO Spheres as Bifunctional Hydrogen Evolution Catalyst and Supercapacitor Electrode

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## Highlights

- MoS<sub>2</sub> nanosheets arrays were vertically rooted on hollow rGO spheres (h-rGO@MoS<sub>2</sub>) via an optimized dual-template strategy.
- The bifunctional h-rGO@MoS<sub>2</sub> architecture exhibit enhanced hydrogen evolution reaction performance (105 mV, onset potential) and higher specific capacitance (238 F g<sup>-1</sup> at 0.5 A g<sup>-1</sup>) as a supercapacitor electrode than pristine MoS<sub>2</sub>.

**Abstract** MoS<sub>2</sub> has attracted attention as a promising hydrogen evolution reaction (HER) catalyst and a supercapacitor electrode material. However, its catalytic activity and capacitive performance are still hindered by its aggregation and poor intrinsic conductivity. Here, hollow rGO sphere-supported ultrathin MoS<sub>2</sub> nanosheet arrays (h-rGO@MoS<sub>2</sub>) are constructed via a dual-template approach and employed as bifunctional HER catalyst and supercapacitor electrode material. Because of the expanded

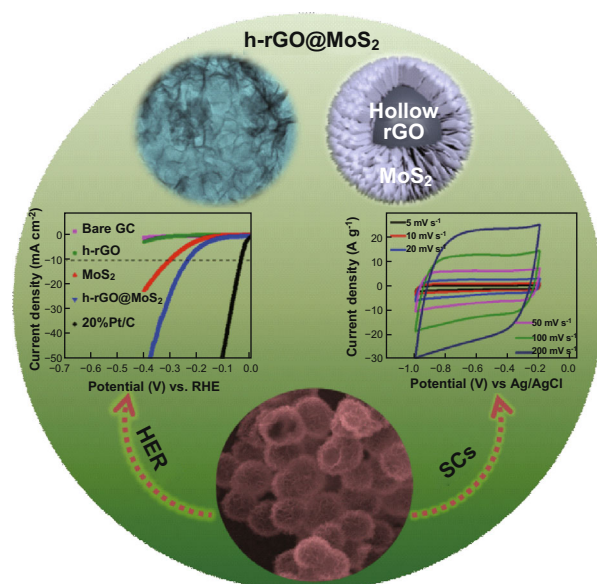
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interlayer spacing in MoS<sub>2</sub> nanosheets and more exposed electroactive S–Mo–S edges, the constructed h-rGO@MoS<sub>2</sub> architectures exhibit enhanced HER performance. Furthermore, benefiting from the synergistic effect of the improved conductivity and boosted specific surface areas (144.9 m<sup>2</sup> g<sup>-1</sup>, ca. 4.6-times that of pristine MoS<sub>2</sub>), the h-rGO@MoS<sub>2</sub> architecture shows a high

specific capacitance ( $238 \text{ F g}^{-1}$  at a current density of  $0.5 \text{ A g}^{-1}$ ), excellent rate capacitance, and remarkable cycle stability. Our synthesis method may be extended to construct other vertically aligned hollow architectures, which may serve both as efficient HER catalysts and supercapacitor electrodes.

**Keywords** MoS<sub>2</sub> · Reduced graphene oxide (rGO) · Hollow spheres · Hydrogen evolution reaction (HER) · Supercapacitor

## 1 Introduction

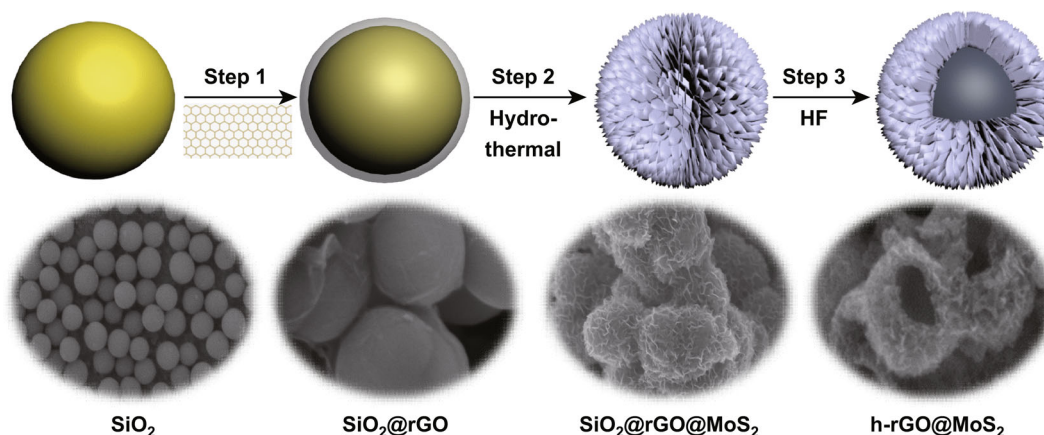
The worsening energy crisis and environmental pollution have stimulated increased research into exploiting sustainable, renewable energy sources, and advanced energy-storage devices. Hydrogen, a clean energy source with the highest gravimetric energy density ( $143 \text{ kJ g}^{-1}$ ) [1], is considered as a promising alternative to fossil fuels and has, thus, attracted significant attention. Water electrolysis is a simple way to produce highly pure hydrogen; but the best-performed Pt-based electrocatalysts suffer from their high cost and scarcity of platinum [2, 3]. Additionally, due to the low cost, rapid charge-discharge process, and long cycling stability, supercapacitors have emerged as promising energy-storage devices to meet the burgeoning demand [4, 5]. Nevertheless, the energy density of supercapacitors is still less than satisfactory [6].

To address this bottleneck, emerging 2D materials, particularly MoS<sub>2</sub>, have been widely studied as hydrogen evolution reaction (HER) catalysts [7–14] and supercapacitor electrode materials [15, 16]. Unfortunately, because of the poor intrinsic conductivity and easy aggregation, the HER catalytic activity and supercapacitive performance of bulk MoS<sub>2</sub> are poor [17, 18]. Accordingly, significant efforts have been devoted to confine the growth of MoS<sub>2</sub> using conductive matrix as template [19, 20]. Reduced graphene oxide (rGO), a layered carbon material with large-specific surface area and excellent electrical conductivity, has been proven to be an effective matrix that can endow MoS<sub>2</sub> with specific shapes, expand the interlayer spacing, and increase the conductivity [17, 21–26]. For example, Dai and coworkers synthesized a MoS<sub>2</sub>/rGO hybrid structure via a solvothermal process, and the product benefitted from the synergistic effect of strong chemical and electronic coupling effects, an abundance of electroactive edges, and improved conductivity; thus, the hybrid exhibited enhanced HER performance [17].

Although remarkable progress has been achieved toward improving the HER performance by introducing graphene oxide as a growth matrix, the performance is still far from satisfactory. On the one hand, the inevitable restacking of graphene oxide, arising from the strong sheet-to-sheet  $\pi$ - $\pi$  interactions [27], hampers the full use of the active surfaces. On the other hand, the limited ion diffusion and mass transport of the flat structure lead to the unsatisfactory HER catalytic activity.

Construction of hierarchical hollow architectures is deemed to be an effective strategy to enhance the electrochemical properties because of the advantages of a high surface-to-volume ratio and open structure [28–30]. In addition, the cavities in the hollow structures can serve as “ion-buffering reservoirs” to shorten the ion transport distances [31]. Previous papers [32, 33] have reported the application of conductive carbon shells as supports, yielding hierarchical hollow architectures; however, the lack of functional groups on surface constrains the growth of MoS<sub>2</sub> to ultrathin nanosheets. Moreover, the relatively thick carbon shells ( $> 20 \text{ nm}$ ) have adverse effect on ion diffusion and mass transport. Based on the above considerations, a facile SiO<sub>2</sub>-template-based method is designed to fabricate hollow graphene spheres as supports to confine the growth of MoS<sub>2</sub> nanosheet arrays. It is anticipated that the restacking of GO would be prevented by SiO<sub>2</sub> templates, and MoS<sub>2</sub> is expected to be vertically supported on the GO shells. This expected hollow architecture possesses advantages of enhanced surface area, as well as more exposed electroactive sites, and, thus, is expected to exhibit enhanced HER catalytic activity and supercapacitive performance [34]. However, there have been only few reports on the use of hollow rGO spheres as a matrix to construct hierarchical architectures that serve both as HER catalysts and supercapacitor electrode materials.

In this study, 3D rGO hollow sphere-supported ultrathin MoS<sub>2</sub> nanosheets were prepared, as illustrated in Scheme 1. First, GO was intimately coated around positively charged SiO<sub>2</sub> spheres by electrostatic incorporation. Subsequently, ultrathin MoS<sub>2</sub> nanosheets were vertically grown around the GO shells via a hydrothermal process. Finally, the SiO<sub>2</sub> cores were fully removed by chemical etching. In such 3D architectures, the vertically aligned MoS<sub>2</sub> nanosheets with a large number of electroactive S–Mo–S edges and expanded (002) interlayer spacings were observed to be tightly anchored on the graphene surface. Benefitting from the expanded (002) interlayer spacing and the exposed electroactive S–Mo–S edges, the h-rGO@MoS<sub>2</sub> structure exhibited efficient HER catalytic activity with an overpotential of ca.  $230 \text{ mV}$  at  $10 \text{ mA cm}^{-2}$ . In addition, because of the boosted specific surface area and conductivity, remarkable supercapacitive



**Scheme 1** Schematic illustration of the synthetic process of h-rGO@MoS<sub>2</sub>

performance with high specific capacitance (238 F g<sup>-1</sup> at 0.5 A g<sup>-1</sup>) was achieved.

## 2 Experimental Section

### 2.1 Synthesis of the h-rGO@MoS<sub>2</sub> Architecture

First, GO sheets and SiO<sub>2</sub> spheres were prepared via a modified Hummers' method [35] and typical Stöber method [36], respectively. Before the GO wrapping process, SiO<sub>2</sub> spheres were decorated with (3-aminopropyl) triethoxysilane (APTES) [37] to obtain a positively charged surface. Second, during the GO coating process, 0.5 g of positively charged SiO<sub>2</sub> was dispersed in 300 mL of DI-water in a supersonic bath, and, then, 50 mL of 0.1 mg L<sup>-1</sup> GO was dropped into the solution with continuous stirring. Subsequently, the SiO<sub>2</sub>@GO powders were obtained after filtration and drying in an oven at 60 °C, followed by grinding. Next, 0.5 g of the as-synthesized SiO<sub>2</sub>@GO powder was re-dispersed in 35 mL DI-water, followed by the addition of 0.3 g of Na<sub>2</sub>MoO<sub>4</sub>·2H<sub>2</sub>O and 0.8 g of thiourea to the solution. Subsequently, 1 mL concentrated HCl (36 wt%) was added to adjust the pH value. After 1 h stirring, the blue mixture was transferred to a 50-mL Teflon-lined stainless-steel autoclave and kept at 200 °C for 24 h. After that, the black SiO<sub>2</sub>@rGO@MoS<sub>2</sub> sample was washed thoroughly and dried at 60 °C, followed by vacuum calcination at 350 °C for 2 h. Finally, the SiO<sub>2</sub> templates were removed via etching in 10% HF while stirring for 12 h, and h-rGO@MoS<sub>2</sub> hollow spheres were obtained. For comparison, SiO<sub>2</sub>/MoS<sub>2</sub> and pristine MoS<sub>2</sub> were synthesized under the same conditions without the addition of GO and SiO<sub>2</sub>@rGO template. Additionally, hollow rGO(h-rGO) spheres were

synthesized via GO coating, hydrothermal reduction, and a SiO<sub>2</sub> etching process.

### 2.2 Characterization

The as-prepared samples were characterized by X-ray diffraction (XRD) using a Cu K $\alpha$  radiation source ( $\lambda = 1.5406 \text{ \AA}$ ) operating at 40 kV and 100 mA. Field-emission scanning electron microscopy (FE-SEM, JEOL-6701F), transmission electron microscopy (TEM, JEOL-2010), and high-resolution TEM (HRTEM, JEOL-2010) measurements were also carried out. X-ray photoelectron spectroscopy (XPS) measurements were carried out using a Thermo Scientific ESCALAB 250Xi instrument. The specific surface areas were analyzed using a Bei Shi De (3H-2000PM2) instrument and calculated using the Brunauer–Emmett–Teller (BET) model. The Raman spectra were recorded using an RTS-HiR-AM with excitation at 532 nm and a power of 5  $\mu$ W.

### 2.3 Electrochemical Evaluations

For HER evaluation, the electrochemical tests were carried out in a three-electrode system with 0.5 M H<sub>2</sub>SO<sub>4</sub> as the electrolyte and a Ag/AgCl electrode and graphite rod as the reference and counter electrodes, respectively. For these experiments, 10 mg pristine MoS<sub>2</sub> and 10 mg h-rGO@MoS<sub>2</sub> were dispersed in 1 mL liquor (0.75 mL water + 0.15 mL alcohol + 0.1 mL 5 wt% Nafion) and sonicated for 2 h. Then, 8  $\mu$ L liquor was dropped onto the glassy carbon (GC) electrode using a microsyringe and dried at room temperature, and the mass loading of catalyst was calculated to be 400  $\mu$ g cm<sup>-2</sup>. The electrolytes were bubbled with N<sub>2</sub> before measurement and flowed over the electrolyte during the scanning process to exclude O<sub>2</sub>. The measured potentials were referenced to the reversible

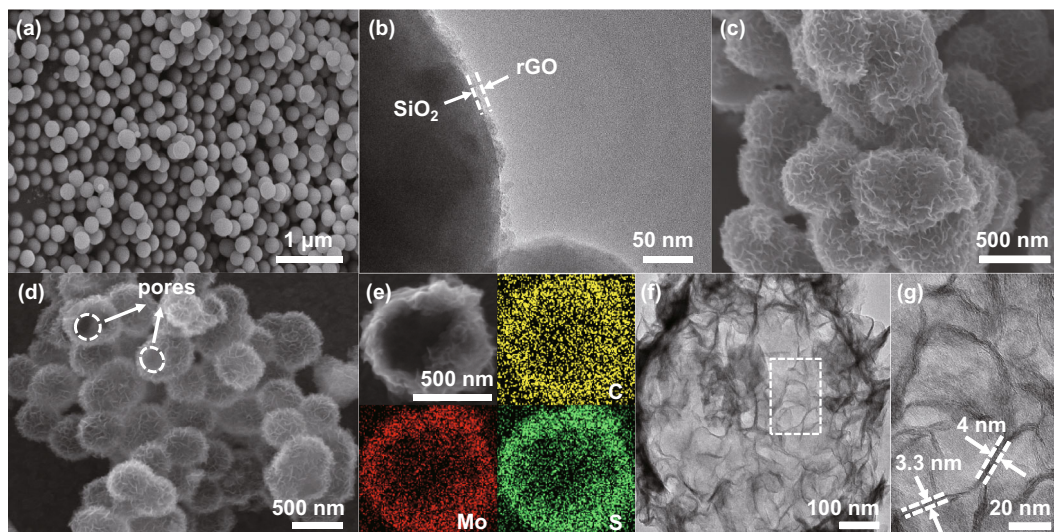
hydrogen electrode (RHE) by adding a value of  $(0.197 + 0.059 \text{ pH}) \text{ V}$ . All data are presented without  $iR$  compensation.

For the evaluation of the supercapacitor, the working electrodes were prepared similarly to previous methods [38]. First, as-prepared samples ( $\text{MoS}_2$  and  $\text{h-rGO@MoS}_2$ ), carbon black, and polyvinylidene difluoride (PVDF) were mixed in mass ratio of 8:1:1 and ground to slurry by adding *N*-methyl pyrrolidone (NMP). Second, the obtained slurry was then painted onto carbon cloth current collectors, yielding a calculated active area of  $1 \text{ cm}^2$ . Finally, the resultant working electrodes were dried at  $60 \text{ }^\circ\text{C}$  for 24 h and pressed at 10 MPa. All electrochemical experiments were performed in 1 M  $\text{Na}_2\text{SO}_4$  solution with platinum foils as the counter electrode and  $\text{Ag/AgCl}$  as the reference electrode, respectively. An electrochemical work station (CH Instruments, CHI760e) was employed for the cyclic voltammetry (CV), electrochemical impedance spectroscopy (EIS) measurements, and galvanostatic charge–discharge (GCD) tests. The EIS measurements were recorded with an AC voltage amplitude of 5 mV and a frequency range from 100 kHz to 0.01 Hz. Galvanostatic charge–discharge investigations were carried out from  $-1$  to  $-0.2 \text{ V}$  at current densities of 0.5, 1, 2, 3, 4, and  $6 \text{ A g}^{-1}$ .

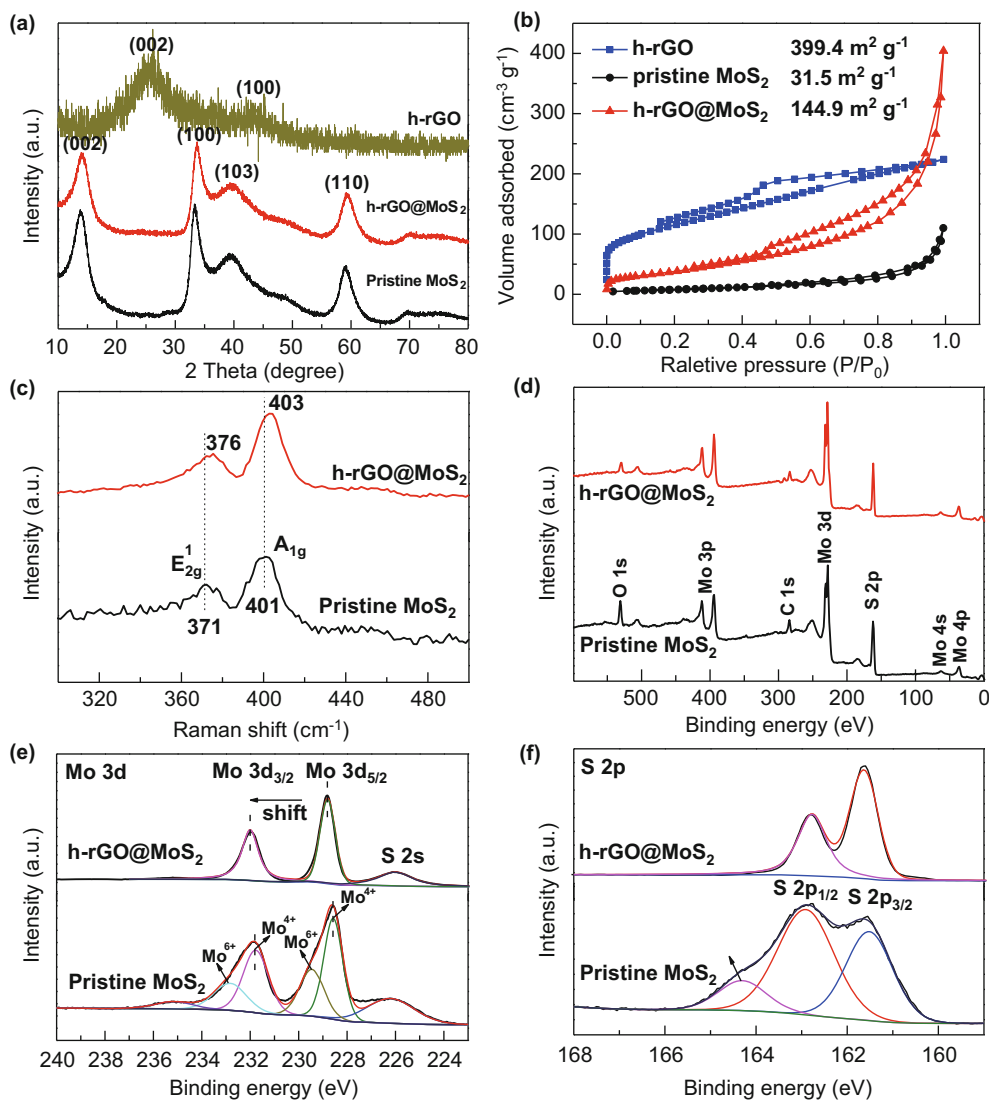
### 3 Results and Discussion

To verify the morphologies of the  $\text{SiO}_2$ ,  $\text{SiO}_2@\text{GO}$ , and  $\text{h-rGO@MoS}_2$  architecture, SEM and TEM were used. Figure 1a shows an SEM image of the  $\text{SiO}_2$  template prepared via the Stöber method, in which the  $\text{SiO}_2$  particles

are uniform spheres with diameters of ca. 350 nm. After decoration with APTES and coating with GO, the glossy surface of  $\text{SiO}_2$  was wholly wrapped with GO, the thickness of which is less than 10 nm (Fig. 1b). Accordingly, because of the template effect of the  $\text{SiO}_2$  spheres, the aggregation of GO sheets was hindered, and the available surface area for  $\text{MoS}_2$  growth was increased. As shown in Fig. 1c, the thin  $\text{MoS}_2$  nanosheets are vertically rooted on  $\text{SiO}_2@\text{GO}$  surface, forming a unique 3D hierarchical architectures. This arises from the abundant hydrophilic functional groups anchored on the surface of GO, that can attract and adsorb  $\text{MoO}_4^{2-}$  [39], resulting in the formation of well-dispersed  $\text{MoS}_2$  nanosheets during the hydrothermal process. Notably, none of scattered  $\text{MoS}_2$  nanosheets can be seen in the SEM image, revealing the appropriate addition of Mo precursors. As a comparison,  $\text{SiO}_2$  spheres without a GO layer were used directly as the growth template, and no  $\text{MoS}_2$  nanosheets were found at the roots (Fig. S1), which can be ascribed to the lack of anchoring nucleation sites for  $\text{MoO}_4^{2-}$  on the  $\text{SiO}_2$  surface. Moreover, because of the stacking nature, pristine  $\text{MoS}_2$  sheets synthesized without  $\text{SiO}_2@\text{GO}$  templates tended to tangle and aggregate, finally forming stacked “solid flowers” (Fig. S2). Therefore, the GO in this work serves as a template to confine the growth of the  $\text{MoS}_2$  nanosheets and inhibit aggregation. To achieve a hollow architecture, HF was used as the etching agent to remove the  $\text{SiO}_2$  template, and the final  $\text{h-rGO@MoS}_2$  spheres are shown in Fig. 1d. As expected,  $\text{h-rGO@MoS}_2$  retains its initial 3D architecture with a diameter of about 480 nm. Due to the 3D architecture,  $\text{h-rGO@MoS}_2$  has a large surface area and more exposed S–Mo–S electroactive edges, which was



**Fig. 1** The SEM and TEM images of **a** the  $\text{SiO}_2$  sphere templates, **b**  $\text{SiO}_2@\text{GO}$ , and **c**  $\text{SiO}_2@\text{GO@MoS}_2$ . **d** SEM images of the  $\text{h-rGO@MoS}_2$  spheres. **e** EDS elemental mapping results of the  $\text{h-rGO@MoS}_2$  spheres. C arises from rGO, Mo and S arises from  $\text{MoS}_2$ . **f** TEM image of  $\text{h-rGO@MoS}_2$ , and **g** the HRTEM image of the white area in **f** showing several sheets of  $\text{MoS}_2$



**Fig. 2** **a** XRD spectra of h-rGO, pristine MoS<sub>2</sub>, and h-rGO@MoS<sub>2</sub>. **b** N<sub>2</sub> adsorption–desorption curves of h-rGO, pristine MoS<sub>2</sub>, and h-rGO@MoS<sub>2</sub>. The corresponding specific surface areas are given in the top left corner. **c** Raman spectra of pristine MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub> measured with 532 nm excitation. **d** Survey spectra of pristine MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub>. **e**, **f** The high-resolution Mo 3d and S 2p XPS spectra of pristine MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub>

confirmed by N<sub>2</sub> adsorption–desorption experiments (Fig. 2b).

To reveal the elemental distribution, energy-dispersive X-ray spectroscopy (EDS) analysis of the h-rGO@MoS<sub>2</sub> was conducted, as shown in Fig. 1e. As shown, a hollow sphere with a cavity is observed, confirming the formation of a hollow architecture. Moreover, elemental C is well spatially distributed at the core, whereas Mo and S are scattered on the outer surface of the shells, indicating that MoS<sub>2</sub> is well dispersed on the surface of rGO. To obtain further information about the h-rGO@MoS<sub>2</sub> architecture, TEM observation was conducted, as shown in Fig. 1f. As clearly shown in the image, 3D hierarchical hollow spheres with cavities were formed, further revealing the successful

formation of hollow architecture. Notably, the existence of these cavities can work as “ion-buffering reservoirs” to reduce the diffusion distance of electrolyte ions [31]. To ensure the layer structure of h-rGO@MoS<sub>2</sub>, corresponding HRTEM measurements are performed in Fig. 1g. As can be seen, the vertically aligned MoS<sub>2</sub> nanosheets have thicknesses of about 3.3–4 nm (5–6 layers), thinner than that of pristine MoS<sub>2</sub> (Fig. S2). Moreover, the interlayer spacing of the MoS<sub>2</sub> nanosheets was measured to be ca. 0.65 nm (Fig. S3), which is assigned to the expanded (002) interlayer spacing of hexagonal MoS<sub>2</sub> [40] and agrees well with the results of XRD analysis (Fig. S4). The expanded interlayer spacing can increase the diffusion kinetics of ions and the interlayer conductivity [41] and, thus, could

enhance the HER performance. As a result, 3D hollow rGO spheres confine the growth of MoS<sub>2</sub> to form highly dispersed and vertically aligned ultrathin nanosheets with large surface area, exposed S–Mo–S electroactive edges, and expanded (002) interlayer spacing, which promise remarkable HER catalytic performance and supercapacitance.

The crystallographic structures (Fig. 2a) of h-rGO, pristine MoS<sub>2</sub>, and h-rGO@MoS<sub>2</sub> architecture were identified via XRD characterization. In the XRD analysis of rGO, two broad peaks appear at  $2\theta = 25^\circ$  and  $43^\circ$ , demonstrating that graphene oxide was reduced to rGO. With respect to the XRD spectra of MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub>, both patterns contain peaks corresponding to semi-conductive hexagonal crystalline (2H) MoS<sub>2</sub>, confirming that no impurity phases were introduced via the SiO<sub>2</sub> template method. It is worth noting that the (002) diffraction peak in Fig. S4 is shifted from  $14.1^\circ$  (MoS<sub>2</sub>) to  $13.7^\circ$  in h-rGO@MoS<sub>2</sub>, further revealing the expansion of (002) interlayer spacing. The broad diffraction peak corresponding to amorphous SiO<sub>2</sub> (located at  $2\theta = 22.5^\circ$ , Fig. S5) was not detected, confirming that the SiO<sub>2</sub> templates had been successfully removed. In addition, no diffraction peaks of rGO were present on account of its low weight ratio.

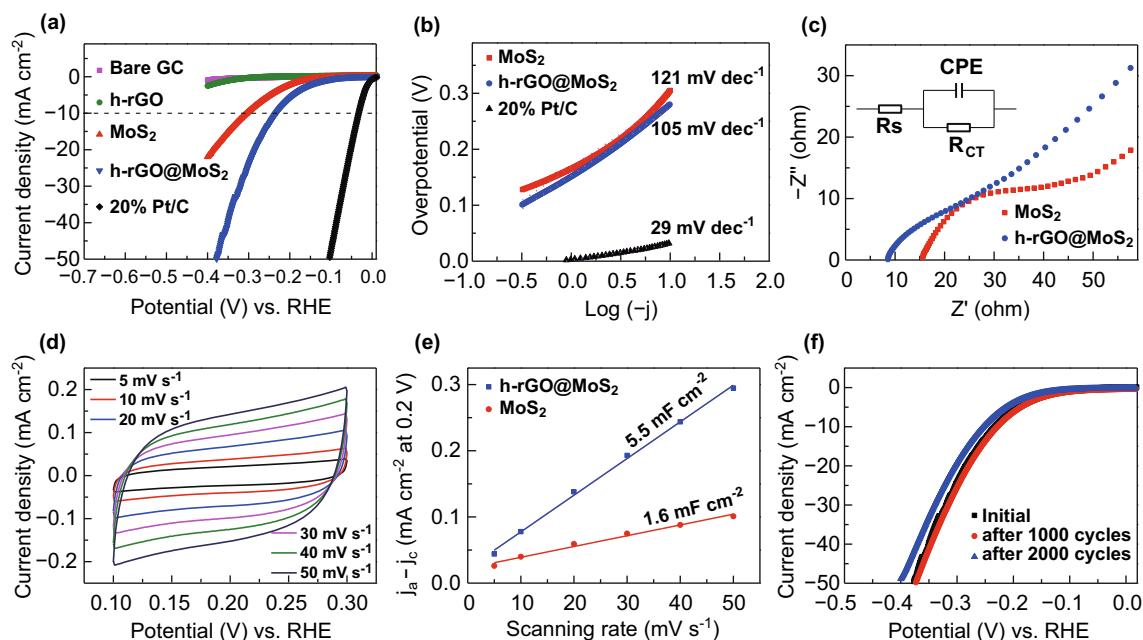
The specific surface area is a key factor that has a significant influence on the HER catalytic performance and specific capacitance. To evaluate the specific surface area, N<sub>2</sub> adsorption–desorption characterization was carried out for h-rGO, pristine MoS<sub>2</sub>, and the h-rGO@MoS<sub>2</sub> architecture, and the results are shown in Fig. 2b. All the curves present type-IV hysteresis loops (IUPAC classification), which may be closely related to the flower-like MoS<sub>2</sub> or 3D hollow rGO and h-rGO@MoS<sub>2</sub> architectures, in accordance with the SEM observations. Because of the hollow structure and unstacked layers, h-rGO spheres exhibit the highest specific surface area of about  $399.4 \text{ m}^2 \text{ g}^{-1}$ . Moreover, after the MoS<sub>2</sub> nanosheets had vertically grown on the surface of h-rGO, the BET surface area of the h-rGO@MoS<sub>2</sub> architecture was calculated to be  $144.9 \text{ m}^2 \text{ g}^{-1}$ , which is about 4.6-times that of pristine MoS<sub>2</sub> ( $31.5 \text{ m}^2 \text{ g}^{-1}$ ). This boosted surface area originates from the confined growth of MoS<sub>2</sub> and vertically aligned nanosheets, as shown in the SEM image of Fig. 1d. Moreover, the confined growth of MoS<sub>2</sub> indirectly results in the exposure of electroactive S–Mo–S edges [12]. This boosted surface area combined with more exposed electroactive S–Mo–S edges may contribute to the superior HER catalytic activity and enhanced supercapacitive performance.

Raman spectra were recorded to obtain structural information about the pristine MoS<sub>2</sub> and the h-rGO@MoS<sub>2</sub> architecture, as shown in Figs. 2c and S6. For pristine

MoS<sub>2</sub>, the two peaks located at  $371$  and  $401 \text{ cm}^{-1}$  are ascribed to the E<sub>2g</sub><sup>1</sup> and A<sub>g</sub><sup>1</sup> modes, respectively [42], and the frequency difference between the E<sub>2g</sub><sup>1</sup> and A<sub>g</sub><sup>1</sup> vibrational modes is about  $30 \text{ cm}^{-1}$ . Importantly, a redshift is observed after MoS<sub>2</sub> nanosheets had vertically grown on h-rGO, and the frequency difference between the E<sub>2g</sub><sup>1</sup> and A<sub>g</sub><sup>1</sup> vibration modes is reduced to  $27 \text{ cm}^{-1}$ , demonstrating reduced number of MoS<sub>2</sub> layers [40]. This agrees well with the results obtained from the TEM images (Figs. 1g and S2). In addition, the two weak peaks situated at  $1350$  and  $1590 \text{ cm}^{-1}$  correspond to the D and G peaks of h-rGO, respectively [17], further revealing the successful formation of h-rGO@MoS<sub>2</sub> composite with a trace amount of h-rGO.

To identify the chemical composition and surface electronic states of pristine MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub>, XPS measurements were carried out. The survey spectra of pristine MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub> are shown in Fig. 2d and contain peaks corresponding to sulfur, molybdenum, carbon, and oxygen, and the calculated Mo/S atomic ratio is about 1/2, confirming the formation of MoS<sub>2</sub>. As shown in Fig. S7, the C 1s peaks of pristine MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub> can be deconvoluted into three peaks and assigned to the sp<sup>2</sup> hybrid C (C–C/C=C,  $284.8 \text{ eV}$ ), hydroxyl C (C–O,  $286.5 \text{ eV}$ ), and epoxy C (C=O,  $288.5 \text{ eV}$ ), respectively [42]. It should be noted that the carbon in pristine MoS<sub>2</sub> is mainly derived from the carbon additive used during measurement. Importantly, for the C 1s peaks of h-rGO@MoS<sub>2</sub>, the oxygen-containing groups yield low-intensity peaks, demonstrating that GO has been successfully reduced to rGO. The high-resolution Mo 3d spectrum of pristine MoS<sub>2</sub> (Fig. 2e) can be deconvoluted into two doublet peaks located at  $232.8$  and  $229.4 \text{ eV}$  and at  $231.8$  and  $228.6 \text{ eV}$ , assigned to Mo<sup>6+</sup> and Mo<sup>4+</sup>, respectively. In addition, the small peak at  $235.5 \text{ eV}$  probably originated from (NH<sub>4</sub>)<sub>2</sub>MoS<sub>4</sub> formed during the hydrothermal reaction. Concerning the Mo 3d spectrum of h-rGO@MoS<sub>2</sub>, a single doublet peak situated at  $233$  and  $229.8 \text{ eV}$  can be attributed to Mo 3d<sub>3/2</sub> and Mo 3d<sub>5/2</sub>, respectively, revealing the dominance of Mo<sup>4+</sup> in MoS<sub>2</sub>. Notably, a positive shift in the Mo 3d spectrum is observed when forming h-rGO@MoS<sub>2</sub>, demonstrating the intense electron coupling effects between h-rGO and MoS<sub>2</sub> [17]. Likewise, the high-resolution S 2p spectra of pristine MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub> were measured, as shown in Fig. 2f. Two peaks fitted to S 2p<sub>1/2</sub> and 2p<sub>3/2</sub> appeared in the spectra of both MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub>, verifying the existence of terminal S<sup>2-</sup> ions, which are favorable for HER activity [43]. Moreover, the unexpected peak located at  $164.3 \text{ eV}$  in the spectrum of pristine MoS<sub>2</sub> is derived from (NH<sub>4</sub>)<sub>2</sub>MoS<sub>4</sub>, as discussed above.

Linear sweep voltammetry (LSV) was carried out to understand the HER performance of h-rGO@MoS<sub>2</sub>; bare



**Fig. 3** **a** Polarization curves of bare GC, h-rGO, pristine MoS<sub>2</sub>, h-rGO@MoS<sub>2</sub>, and commercial 20% Pt/C in 0.5 M H<sub>2</sub>SO<sub>4</sub> electrolyte at 5 mV s<sup>-1</sup>, and **b** the corresponding Tafel plots of electrocatalysts in **a**. **c** EIS plots of pristine MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub> from 10<sup>-2</sup> to 10<sup>5</sup> Hz. **d** The electrochemical double-layer capacitances of h-rGO@MoS<sub>2</sub> ranging from 0.1 to 0.3 V at various scan rates (5, 10, 20, 30, 40, and 50 mV s<sup>-1</sup>) in 0.5 M H<sub>2</sub>SO<sub>4</sub>. **e** The capacitive current at 0.20 V as a function of scan rate for pristine MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub> in 0.5 M H<sub>2</sub>SO<sub>4</sub> electrolyte. **f** Polarization curves of the h-rGO@MoS<sub>2</sub> electrocatalysts initially and after 1000 and 2000 cycles in 0.5 M H<sub>2</sub>SO<sub>4</sub> electrolyte

GC, h-rGO, pristine MoS<sub>2</sub>, and commercial Pt/C (20%) were also evaluated as a comparison. The corresponding polarization curves in Fig. 3a were obtained at a sweep rate of 5 mV s<sup>-1</sup> in 0.5 M H<sub>2</sub>SO<sub>4</sub> without *i*R compensation. Bare GC and h-rGO exhibit negligible HER performance, even at overpotential of 400 mV. In contrast, commercial Pt/C (20%) shows the best HER performance with near-zero onset overpotential. As for h-rGO@MoS<sub>2</sub>, the onset potential is found at ca. 105 mV, which is much lower than that of pristine MoS<sub>2</sub> (-162 mV). Further, the overpotential of h-rGO@MoS<sub>2</sub> at 10 mA cm<sup>-2</sup> is about 230 mV, about 76 mV lower than that of pristine MoS<sub>2</sub> (306 mV). To elucidate the HER mechanism further, Tafel slopes are extracted from the linear portions of the Tafel plots (Fig. 3b). Commercial Pt/C (20%) exhibits a Tafel slope as low as 29 mV dec<sup>-1</sup>, as previously reported [44]. For pristine MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub>, the Tafel slopes are calculated to be 121 and 105 mV dec<sup>-1</sup>, respectively, demonstrating that the Volmer–Heyrovsky mechanism is the rate-limiting step [17]. The performance of the h-rGO@MoS<sub>2</sub> architecture is compared to those of other previously reported hollow MoS<sub>2</sub> microspheres or rGO/MoS<sub>2</sub> composites in Table S1. Because of the specific hollow architecture with vertically aligned MoS<sub>2</sub> nanosheets, the h-rGO@MoS<sub>2</sub> shows a superior onset overpotential.

EIS tests were carried out for both MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub> to estimate the internal resistance properties. As shown by Fig. 3c, the charge transfer resistance of h-rGO@MoS<sub>2</sub> is about 15 Ω, which is smaller than its counterpart MoS<sub>2</sub> (28 Ω). This smaller internal resistance stems from the increased conductivity and shortened electron transfer pathways [33], which favor ion permeation and electron transfer, thus improving the HER catalytic performance. Besides the internal resistance, the electrochemical active surface area (ECSA) is another significant factor that has an impact on the HER performance. The corresponding electrochemical double-layer capacitances of h-rGO@MoS<sub>2</sub> and MoS<sub>2</sub> were measured by CV (ranging from 0.1 to 0.3 V) at different scan rates, as shown in Figs. 3d and S8. Moreover, a linear relationship between the scan rate and current density is observed, as shown in Fig. 3e. Based on the calculation, h-rGO@MoS<sub>2</sub> has a capacitance of 5.5 mF cm<sup>-2</sup>, which is 3.5 times higher than that of pristine MoS<sub>2</sub>. In other words, the electrochemically active sites are boosted after the vertical growth of MoS<sub>2</sub> on h-rGO, and this increase in electrochemically active sites no doubt enhanced the intrinsic HER activity of MoS<sub>2</sub>. Additionally, long-term cycling tests of h-rGO@MoS<sub>2</sub> are examined by continuous cyclic voltammetry measurements between -0.4 and 0.1 V versus RHE at 50 mV s<sup>-1</sup> in Fig. 3f. Clearly, negligible performance decay is observed after the first 1000 CV

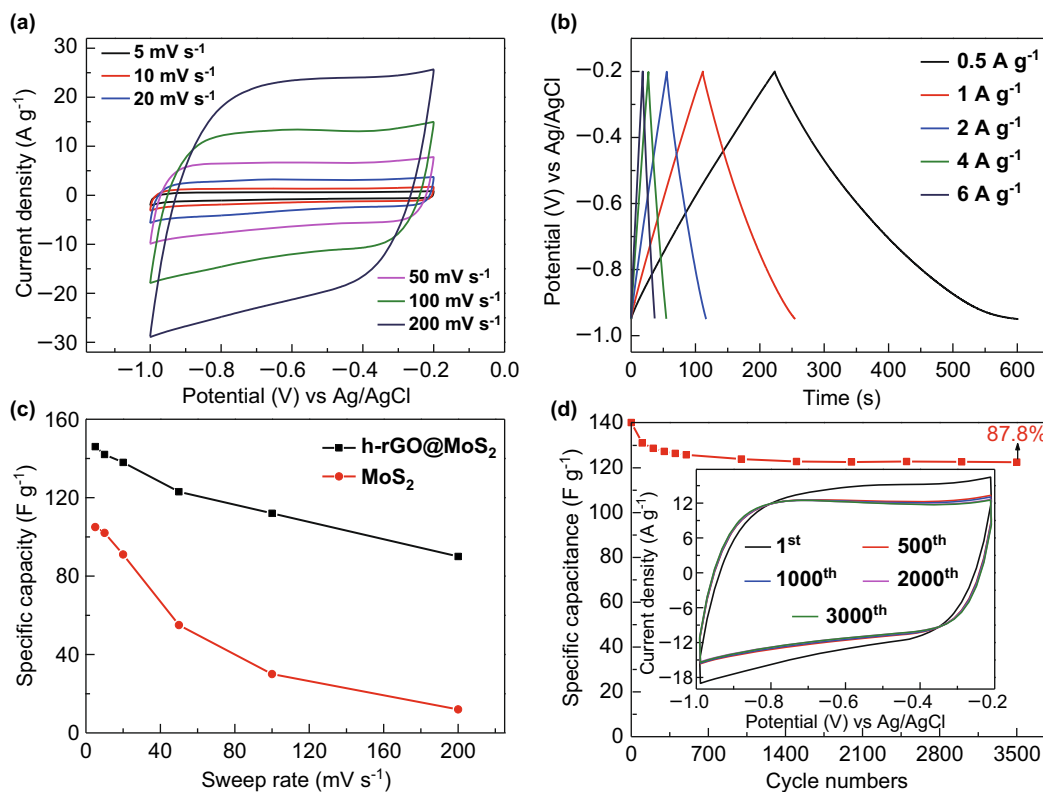
cycles. Moreover, slight decay is seen after 2000 CV cycles, demonstrating the remarkable stability of h-rGO@MoS<sub>2</sub>.

In this study, hollow rGO spheres were used as a matrix to confine the growth of MoS<sub>2</sub> and prevent its agglomeration. Note that the improved HER catalytic activity is derived from a synergistic effect, combining the following three aspects: (1) the expanded (002) interlayer spacing that increases the diffusion kinetics of ions and the intralayer conductivity [41]; (2) the ultrathin vertically aligned MoS<sub>2</sub> with a large number of exposed electroactive sites and increased contact area with the electrolyte [24]; and (3) the improved conductivity that accelerates electron transfer [45]. Thanks to the aforementioned synergistic effects, h-rGO@MoS<sub>2</sub> shows superior HER performance.

Further, the supercapacitive performance of h-rGO@MoS<sub>2</sub> was evaluated using CV and GCD measurements in 1 M Na<sub>2</sub>SO<sub>4</sub> electrolyte. As shown in Fig. 4a, various applied scan rates (ranging from 5 to 200 mV s<sup>-1</sup>) were selected to evaluate the rate performance of h-rGO@MoS<sub>2</sub>. With increasing scan rate, the current densities gradually increased, and they remain almost near-rectangular at all tested scan rates, suggesting the high rate performance of h-rGO@MoS<sub>2</sub>. The specific capacitance is

calculated as  $C_s = \int I dv / v m \Delta V$ , where  $I$  (A),  $v$  (V s<sup>-1</sup>),  $\Delta V$  (V), and  $m$  (g) are the response current, scan rate, potential window, and the mass of active electrode material, respectively. According to the CV curves, the specific capacitances are found to be 146, 142, 138, 123, 112, and 90 F g<sup>-1</sup> at scan rates of 5, 10, 20, 50, 100, and 200 mV s<sup>-1</sup>, respectively. In the case of the pristine MoS<sub>2</sub> electrode, however, the shapes of the CV curves change from near-rectangular to twisted ellipses as the scan rate increased from 5 to 200 mV s<sup>-1</sup>, revealing its less capacitive but more resistive performance (Fig. S9). To compare the capacitance and rate performance of MoS<sub>2</sub> and the h-rGO@MoS<sub>2</sub> architecture, the CV curves are compared at a sweep rate of 100 mV s<sup>-1</sup> in Fig. S10. The CV curve of the h-rGO@MoS<sub>2</sub> architecture has a typical rectangular shape with a large area, while the shape of pristine MoS<sub>2</sub> has become a tilted ellipse.

Regarding the specific capacitance, h-rGO@MoS<sub>2</sub> shows a higher capacitance value (112 F g<sup>-1</sup>) than that of pristine MoS<sub>2</sub> (30 F g<sup>-1</sup>). This high capacitance and good rate performance arise from the formation of vertical 3D architectures, as shown in Fig. 1d, which allow electrolyte ions to access the interior surfaces of the electrode more easily and shorten the permeation distance, resulting in



**Fig. 4** **a** CV curves of h-rGO@MoS<sub>2</sub> electrode at various sweep rates (5, 10, 20, 50, 100, and 200 mV s<sup>-1</sup>) in 1 M Na<sub>2</sub>SO<sub>4</sub> electrolyte. **b** Galvanostatic curves of h-rGO@MoS<sub>2</sub> at various current densities (0.5, 1, 2, 4, and 6 A g<sup>-1</sup>). **c** Capacitance retention property of h-rGO@MoS<sub>2</sub> and MoS<sub>2</sub> at sweep rates from 5 to 200 mV s<sup>-1</sup>. **d** Cycle stability and capacitance retention of the h-rGO@MoS<sub>2</sub> electrode tested at a scan rate of 100 mV s<sup>-1</sup>; the inset shows the CV curves for the 1st, 500th, 1000th, 2000th, and 3000th cycles



increased capacitive performance. Galvanostatic charge–discharge measurements were conducted at various current densities from 0.5 to 6 A g<sup>-1</sup> in a potential window of -1 to -0.2 V. As shown in Fig. 4b, the charge curves are approximately linear and symmetric with their discharge counterparts, which further indicate the double electrode layer capacitance and excellent reversibility of the h-rGO@MoS<sub>2</sub> architecture [16]. The specific capacitances of the h-rGO@MoS<sub>2</sub> architecture are calculated using  $C = I\Delta t/\Delta V$ , where  $C$  is the specific capacitance,  $I$  is the constant discharge current density,  $\Delta t$  is the discharging time, and  $\Delta V$  is the potential window. The specific capacitance of h-rGO@MoS<sub>2</sub> was calculated to be 238 F g<sup>-1</sup> at a current density of 0.5 A g<sup>-1</sup>.

Moreover, even at a high current density of 6 A g<sup>-1</sup> the specific capacitance remained 135 F g<sup>-1</sup>, which demonstrates the excellent rate performance. However, the specific capacitance of pristine MoS<sub>2</sub> can only reach 106 F g<sup>-1</sup> at a current density of 0.5 A g<sup>-1</sup>, and reduced to 41.2 F g<sup>-1</sup> when the current density is increased to 6 A g<sup>-1</sup> (Fig. S11). The supercapacitive performances of MoS<sub>2</sub>/graphene-based materials are summarized and compared in Table S2. As shown, the h-rGO@MoS<sub>2</sub> architecture has a better specific capacitance than the other materials, which probably originates from the synergistic effect of the large surface area for the adsorption–desorption of ions and increased conductivity.

The capacitance during ultrafast charging/discharging was tested for pristine MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub> (the results are shown in Fig. 4c). At a sweep rate of 5 mV s<sup>-1</sup>, the specific capacitance of h-rGO@MoS<sub>2</sub> is ca. 146 F g<sup>-1</sup>, which is ca. 1.4 times higher than that of pristine MoS<sub>2</sub> (105 F g<sup>-1</sup>). As the sweep rate increased from 5 to 200 mV s<sup>-1</sup>, the specific capacitance of the h-rGO@MoS<sub>2</sub> electrode decreased, but still achieved remarkable capacitance retention of 60% with a 40-fold increase in sweep rate. The specific capacitance of pristine MoS<sub>2</sub> decreased sharply with increasing sweep rate and only 10% capacitance is retained at sweep rate of 200 mV s<sup>-1</sup>. The significant discrepancy in the rate capabilities of the h-rGO@MoS<sub>2</sub> electrode and pristine MoS<sub>2</sub> demonstrate that the introduction of graphene and the formation of vertical architectures contributed to boosting the capacitance retention performance.

The cycling stability is another significant factor that impacts the large-scale application of supercapacitors. The evaluation of the cycling stability of pristine MoS<sub>2</sub> and h-rGO@MoS<sub>2</sub> was investigated by CV cycling at a sweep rate of 100 mV s<sup>-1</sup>. As shown in Fig. S12, pristine MoS<sub>2</sub> displayed stable cyclic behavior. With the 3D vertical architecture, the capacitance of the h-rGO@MoS<sub>2</sub> electrode showed an approximately 10% reduction in the first 400 cycles (Fig. 4d), which probably arises from the

volume change of the active materials during the CV processes. In the subsequent 500–3500 cycles, the capacitance of h-rGO@MoS<sub>2</sub> remained steady at 87.7% of its initial capacitance, demonstrating the excellent long-term cycling stability.

## 4 Conclusions

In summary, 3D rGO hollow sphere-supported ultrathin MoS<sub>2</sub> nanosheets have been prepared, and show enhanced HER catalytic activity and supercapacitive performance compared to pristine MoS<sub>2</sub>. The enhanced HER catalytic activity and supercapacitive performance are mainly derived from the following three factors: (1) the vertically aligned hierarchical architecture, which provides a large surface area for the adsorption, desorption, and diffusion of ions; (2) the relatively thin MoS<sub>2</sub> sheets, which provides more active sites and expanded interlayer spacing; and (3) the increased conductivity. Our strategy for constructing 3D architectures of vertically aligned nanosheets supported on hollow spheres may be applied to 3D composites of other materials and offers the potential for the development of efficient HER catalysts and advanced supercapacitor electrodes.

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